

TWIST

The TRIUMF Weak Interaction

Symmetry Test

Precision Muon Decay at TRIUMF

Nathan Rodning

University of Alberta

TWIST: Universities of Alberta, British Columbia,
Northern British Columbia, Montreal, Saskatchewan;
TRIUMF, Texas A&M, Valporaiso, KIAE - Russia

TWIST - Personnel

TRIUMF

- ❖ **Willy Andersson**
- Yuri Davydov
- Jaap Doornbos
- ❖ **Wayne Faszer**
- Dave Gill
- Peter Gumplinger
- Richard Helmer
- Robert Henderson
- John Macdonald
- Glen Marshall
- Art Olin
- ❖ **David Ottewell**
- ❖ **Robert Openshaw**
- Jean-Michel Poutissou
- Renee Poutissou
- ❖ **Grant Sheffer**
- Hans-Christian Walter
- Dennis Wright

Alberta

- ✂ **Andrei Gaponenko**
- Peter Green
- Peter Kitching
- ✂ **Rob MacDonald**
- Maher Quraan
- Nathan Rodning
- ❖ **John Schaapman**
- ✂ **Farhana Sobratee**
- ❖ **Jan Soukup**
- Glen Stinson

British Columbia

- ✂ **Blair Jamieson**
- ❖ **Doug Maas**

- Mike Hasinoff

Northern British Columbia

- Elie Korkmaz
- Tracy Porcelli

Montreal

- Pierre Depommier

✂ **Students**

❖ **Professional Staff**

Regina

- Ted Mathie
- ✂ **George Price**
- Roman Tacik

Saskatchewan

- Bill Shin

Texas A&M

- Carl Gagliardi
- John Hardy
- ✂ **Jim Musser**
- Robert Tribble
- Maxim Vasiliev

Valparaiso

- Don Koetke
- Robert Manweiler
- ❖ **Paul Nord**

- Shirvel Stanislaus

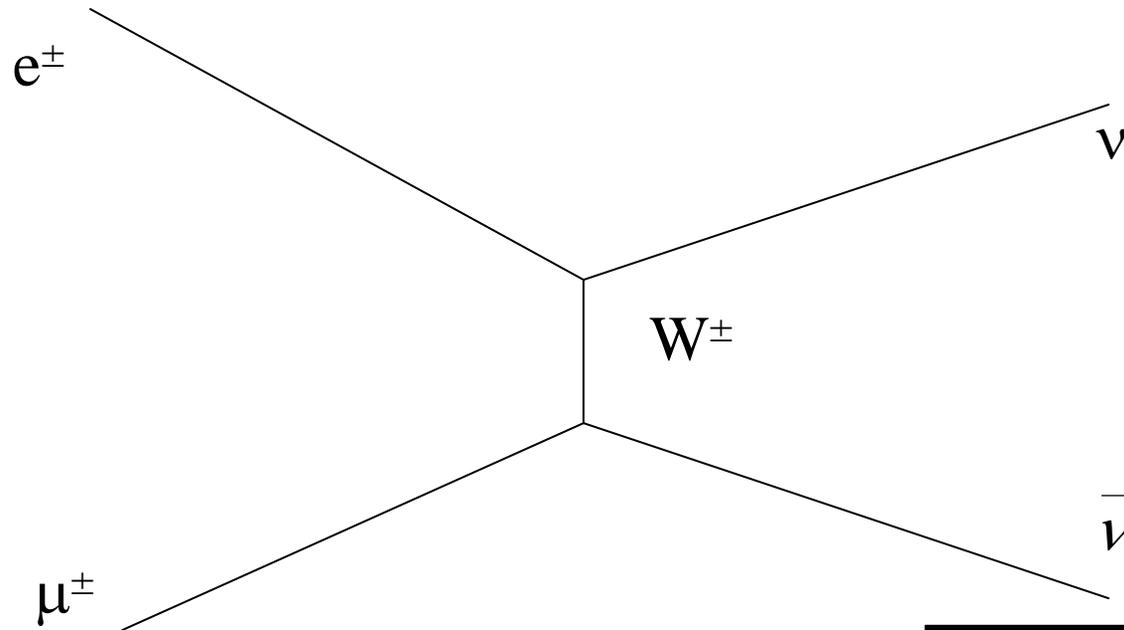
KIAE (Russia)

- Arkadi Khruchinsky
- Vladimir Selivanov
- Vladimir Torokhov

Outline

- **Background on muon decay**
- **The E614 Experiment**
- **Sensitivity to new physics**

The Standard Model for μ decay



(V-A) Interaction is built in

- parity violation is perfect
- exchange particle is known

Only one coupling is non-zero
in the Standard Model

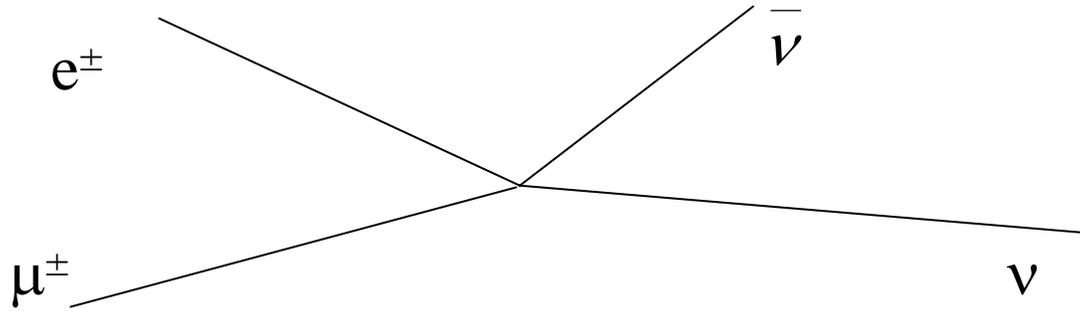
$ g_{RR}^S \equiv 0$	$ g_{RR}^V \equiv 0$	$ g_{RR}^T = \text{zero}$
$ g_{LR}^S \equiv 0$	$ g_{LR}^V \equiv 0$	$ g_{LR}^T \equiv 0$
$ g_{RL}^S \equiv 0$	$ g_{RL}^V \equiv 0$	$ g_{RL}^T \equiv 0$
$ g_{LL}^S \equiv 0$	$ g_{LL}^V \equiv 1$	$ g_{LL}^T = \text{zero}$

- The operator (V-A) satisfies the requirement that the Weak interaction violates parity.
- (V-A) violates parity perfectly
- The (V-A) operator projects out the left-handed (negative chirality) component of the wave function

$$\begin{aligned} \bar{\psi} \gamma^\mu (1 - \gamma^5) \psi &= \bar{\psi} \gamma^\mu (1 - \gamma^5) \begin{bmatrix} \psi_+ \\ \psi_- \end{bmatrix} \\ &= \bar{\psi} \gamma^\mu \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \psi_+ \\ \psi_- \end{bmatrix} = \bar{\psi} \gamma^\mu \psi_- \end{aligned}$$

- the (V-A) theory therefore states that leptons with positive chirality do not undergo weak interactions.

A more general interaction - which does not presuppose the W



$$rate \sim \left| \sum_{\substack{\gamma=S,V,T \\ i,j=R,L}} g_{ij}^\gamma \langle \bar{\psi}_{ei} | \Gamma^\gamma | \psi_{\nu_e} \rangle \langle \bar{\psi}_{\nu_\mu} | \Gamma_\gamma | \psi_{\mu j} \rangle \right|^2$$

Allows for possible

- scalar
- vector
- tensor

Scalar	$\bar{\psi} \psi$
Vector	$\bar{\psi} \gamma^\mu \psi$
Tensor	$\bar{\psi} \sigma^{\mu\nu} \psi$
Axial Vector	$\bar{\psi} \gamma^5 \gamma^\mu \psi$
Pseudoscalar	$\bar{\psi} \gamma^5 \psi$

interactions of right-handed and left-handed leptons

The preceding - in terms of the Michel parameters

$$\text{rate} \sim x^2 \left[3 - 3x + \frac{2}{3} \rho(4x - 3) + P_\mu \xi \cos(\theta) \left(1 - x + \frac{2}{3} \delta(4x - 3) \right) \right]$$

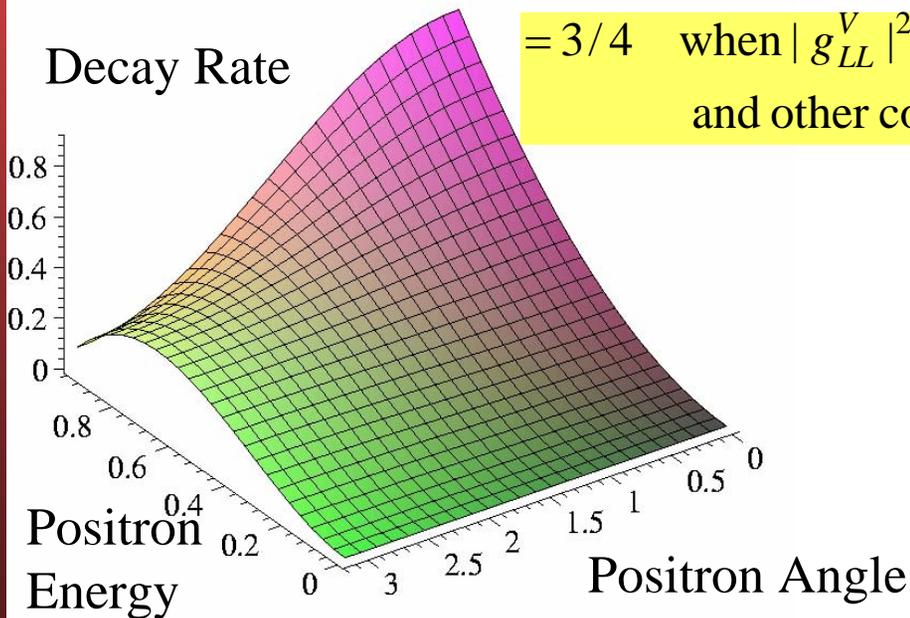
For example-

$$\begin{aligned} \rho \equiv & \frac{3}{4} \left[|g_{LL}^V|^2 + |g_{RR}^V|^2 + |g_{LR}^T|^2 + |g_{RL}^T|^2 \right] \\ & + \frac{3}{16} \left[|g_{LL}^S|^2 + |g_{RR}^S|^2 + |g_{LR}^S|^2 + |g_{RL}^S|^2 \right] \\ & - \frac{3}{4} \left[\text{Re}(g_{LR}^S g_{LR}^{T*}) + \text{Re}(g_{RL}^S g_{RL}^{T*}) \right] \end{aligned}$$

$$= 3/4 \quad \text{when } |g_{LL}^V|^2 = 1 \text{ and other couplings are zero}$$

Similar expressions exist defining ξ and δ .

A fourth parameter, η , contributes to order (m_e/m_μ)



Above expression is modified by radiative corrections, required to second order

The Expression becomes considerably simpler in the Standard Model

$$\text{rate} \sim x^2 \left[3 - 3x + \frac{2}{3} \rho (4x - 3) + P_\mu \xi \cos(\theta) \left(1 - x + \frac{2}{3} \delta (4x - 3) \right) \right]$$

For example-

$$\begin{aligned} \rho \equiv & \frac{3}{4} \left[|g_{LL}^V|^2 + |g_{RR}^V|^2 + |g_{LR}^T|^2 + |g_{RL}^T|^2 \right] \\ & + \frac{3}{16} \left[|g_{LL}^S|^2 + |g_{RR}^S|^2 + |g_{LR}^S|^2 + |g_{RL}^S|^2 \right] \\ & - \frac{3}{4} \left[\text{Re}(g_{LR}^S g_{LR}^{T*}) + \text{Re}(g_{RL}^S g_{RL}^{T*}) \right] \end{aligned}$$

**Standard Model
Values**

$$= 3/4 \quad \text{when } |g_{LL}^V|^2 = 1$$

and other couplings are zero

Similar expressions exist defining ξ , δ , and η .

This simple model may be too simple

exchange particle:

spin 0

spin 1

spin 2

$$|g_{RR}^S| < 0.066$$

$$|g_{RR}^V| < 0.033$$

$$|g_{RR}^T| \equiv 0$$

$$|g_{LR}^S| < 0.125$$

$$|g_{LR}^V| < 0.060$$

$$|g_{LR}^T| < 0.036$$

$$|g_{RL}^S| < 0.424$$

$$|g_{RL}^V| < 0.110$$

$$|g_{RL}^T| < 0.122$$

$$|g_{LL}^S| < 0.55$$

$$|g_{LL}^V| > 0.96$$

$$|g_{LL}^T| \equiv 0$$

All but one of these terms has been set to zero in the Standard model for simplicity

The Weak Interaction may not be purely (V-A)

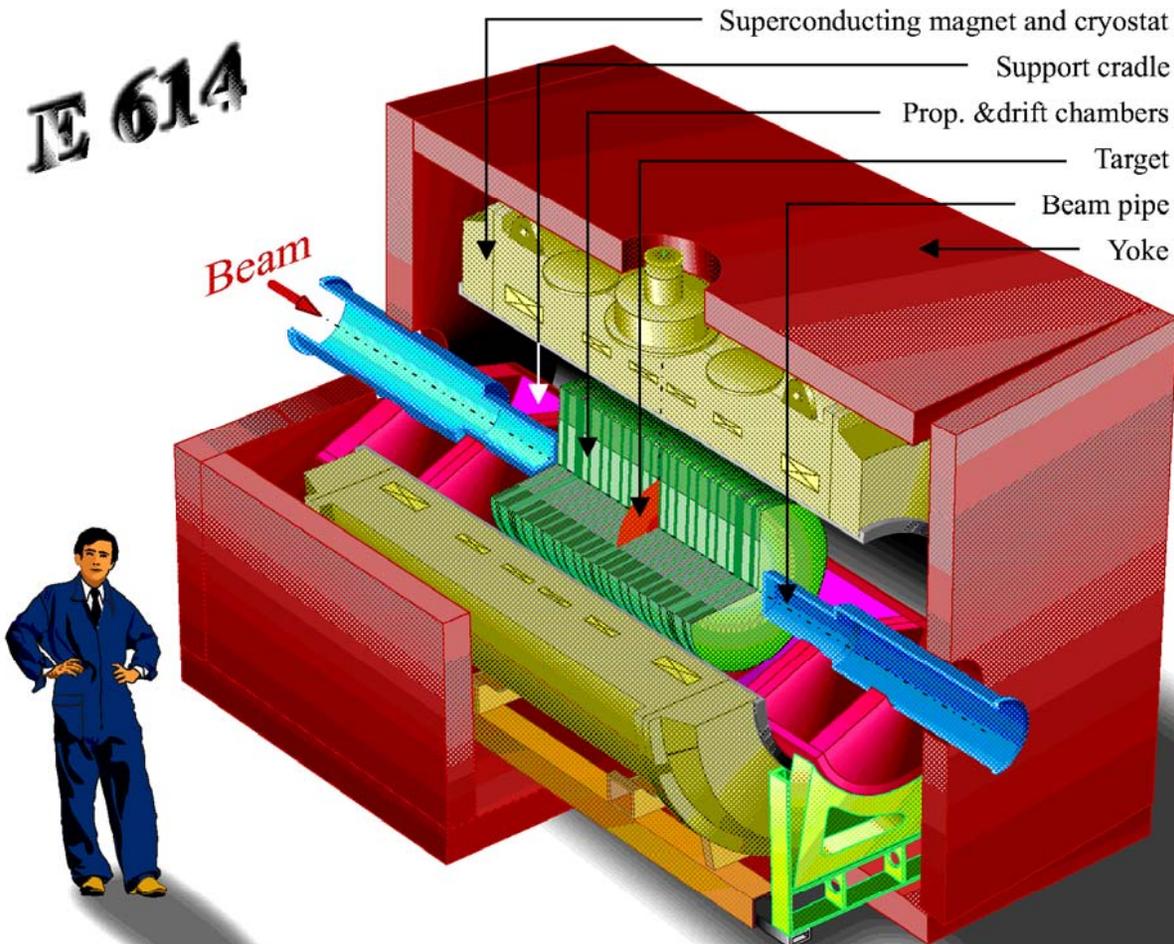
TWIST - Goals

We propose to study 10^9 μ^+ decays

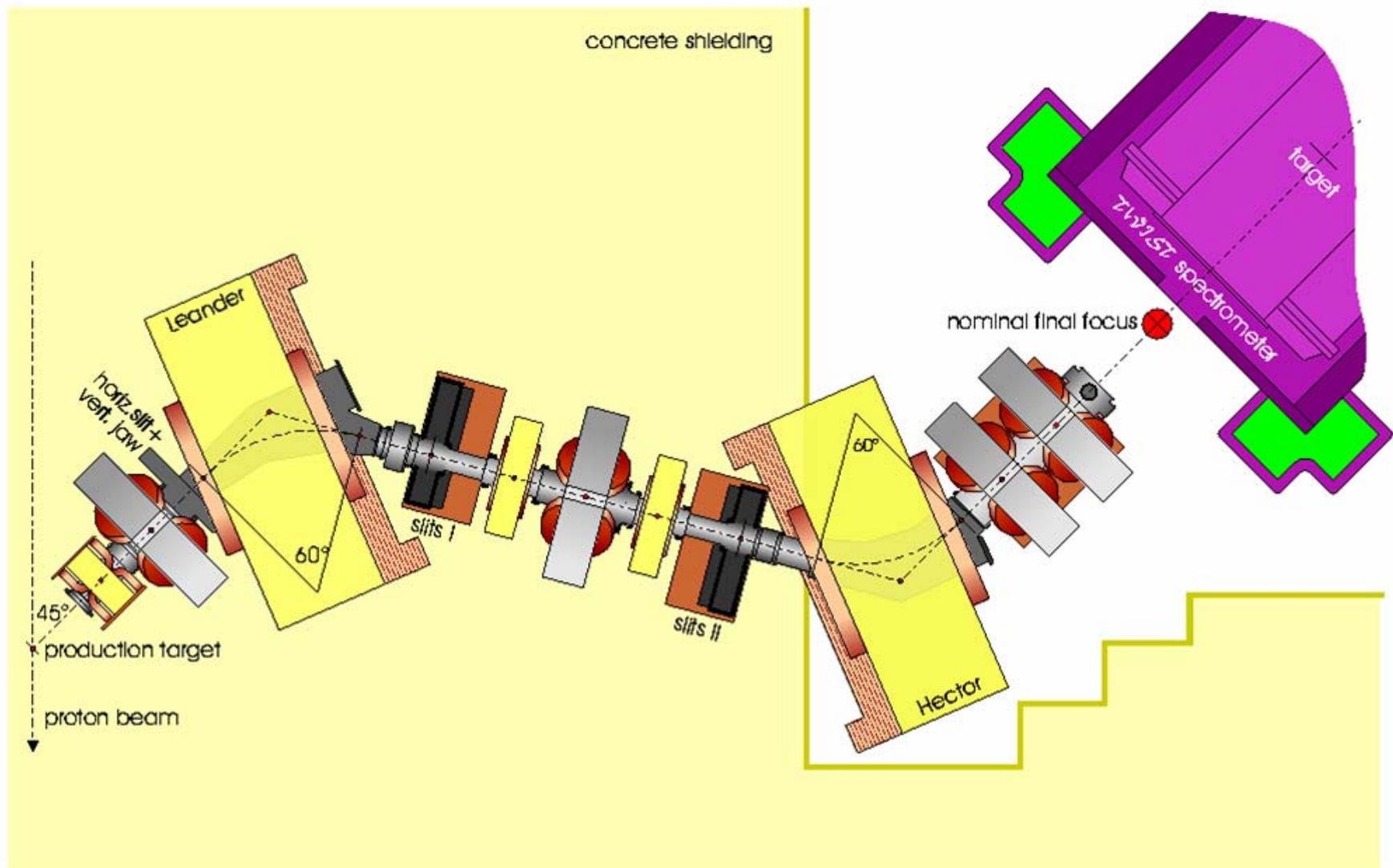
Goal:

- to determine the Michel parameters to a few parts in 10^4
- to test for weak couplings inconsistent with the Standard Model

TWIST - Spectrometer

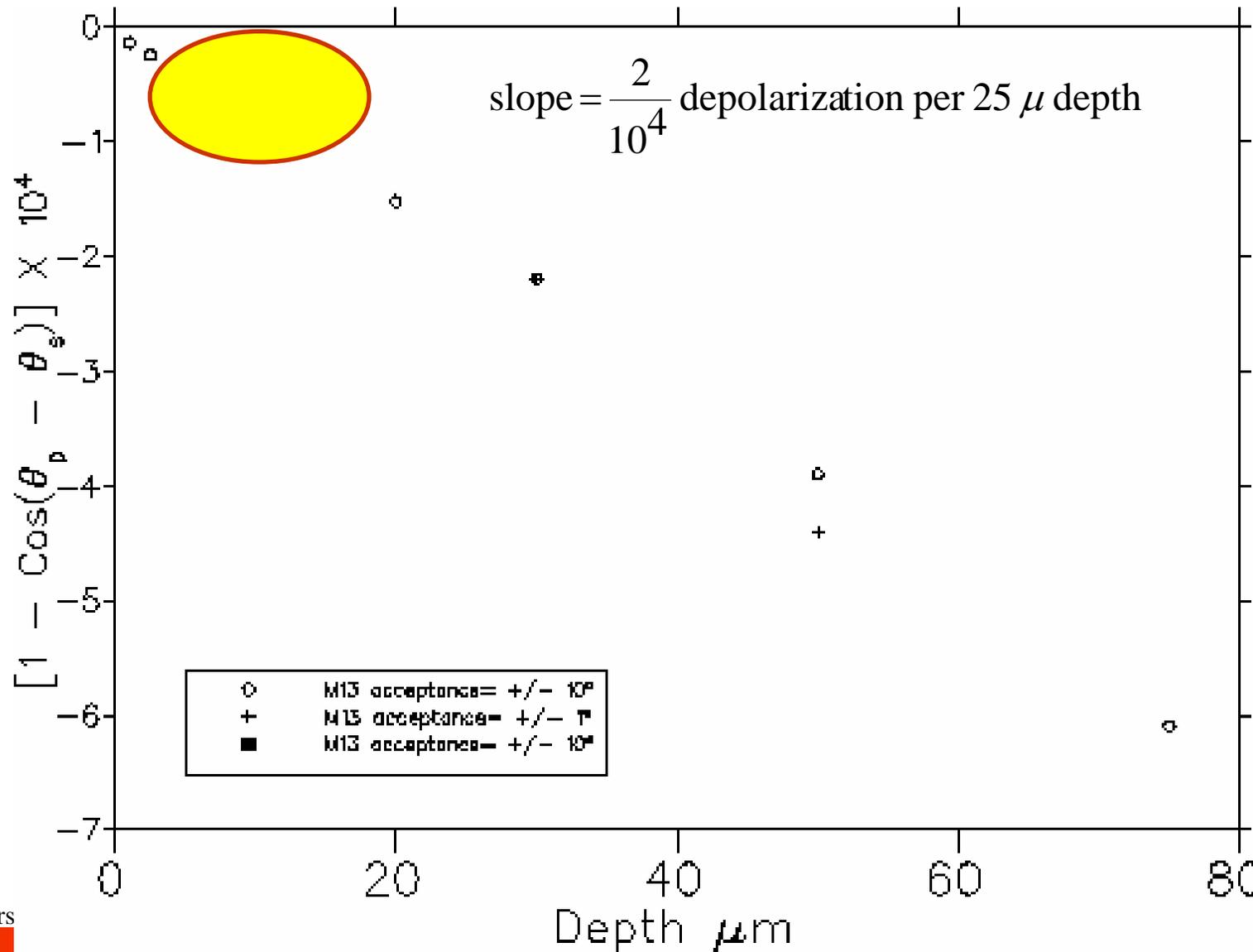


TWIST - Beamline



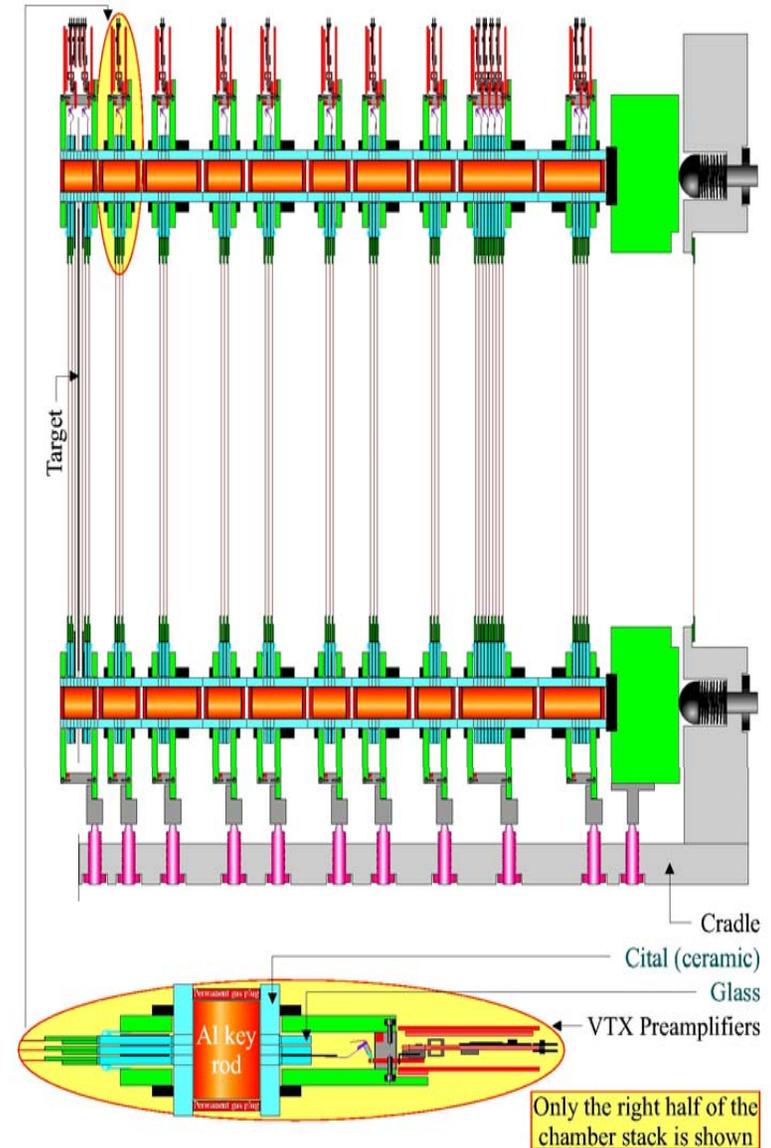
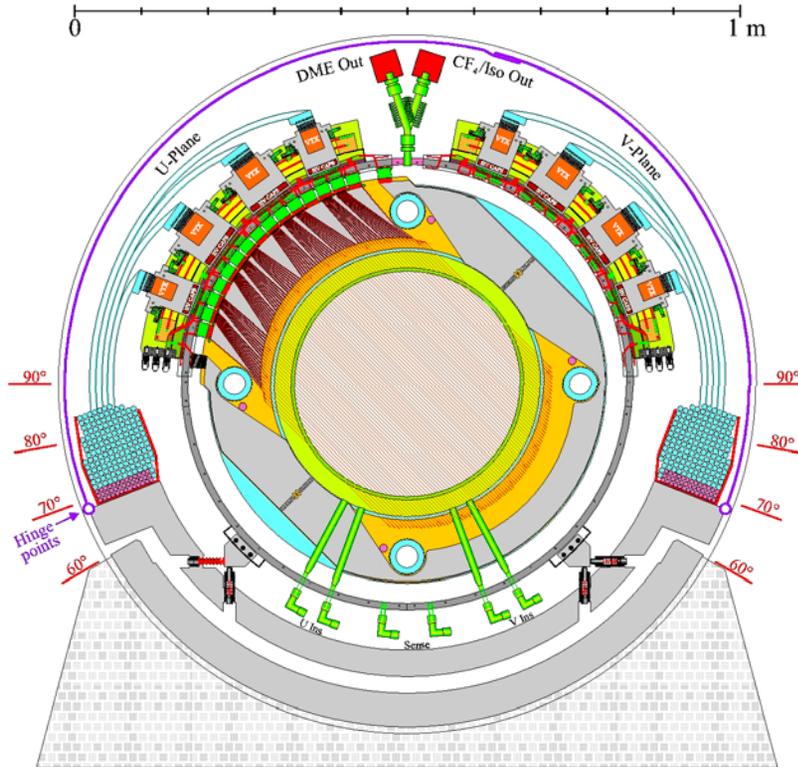
TWIST-1AT1 depolarization

1AT1 Scatter => ~ 0.0001 depolarization



TWIST - Chambers & half detector

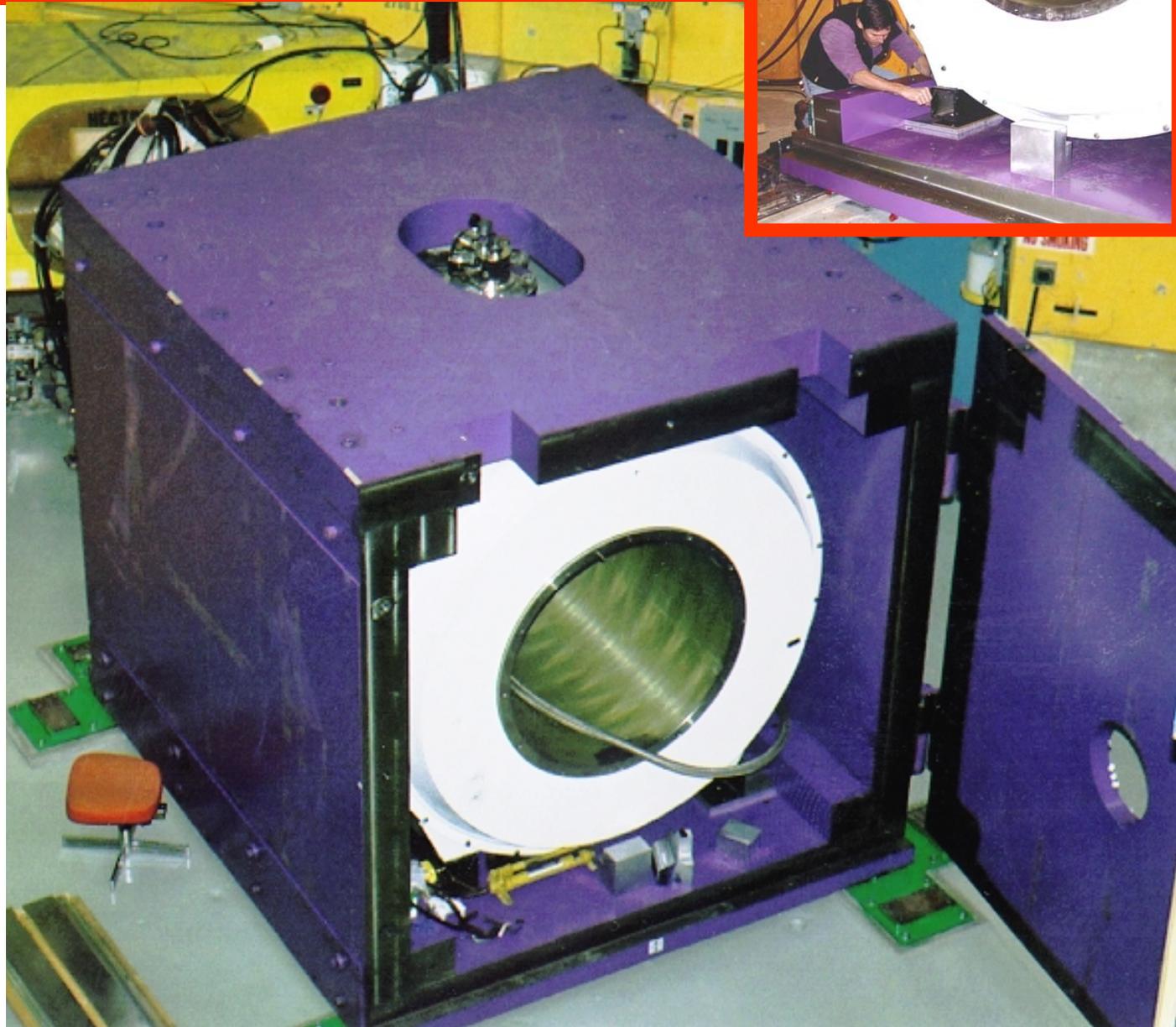
Planar drift chambers sample positron track



TWIST - Yoke

The TWIST yoke pieces were delivered and assembled before Christmas

Alignment was completed in the first week of January



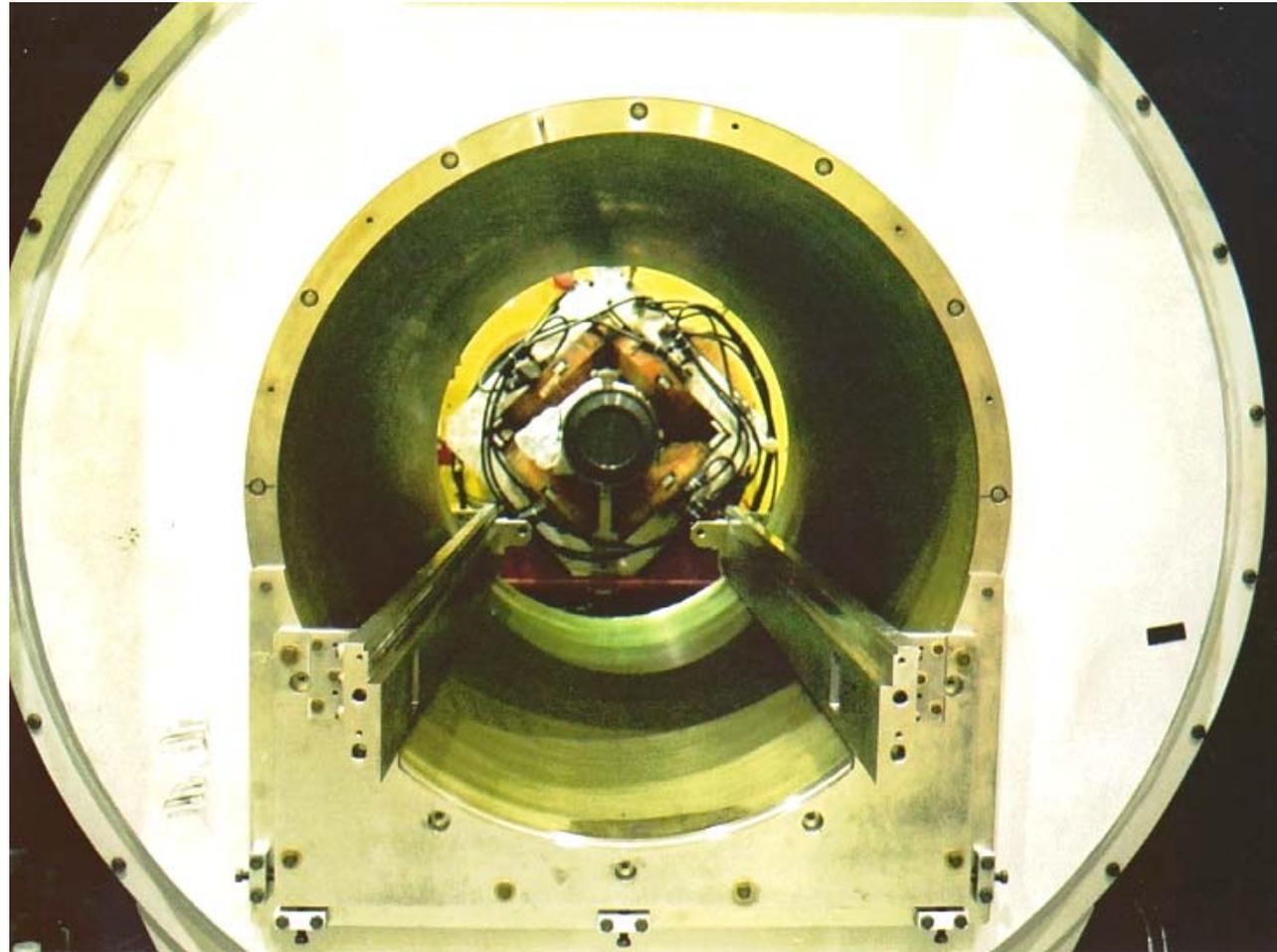
TWIST- Solenoid and WC track

Track is in place and aligned to accept detector cradle and stack

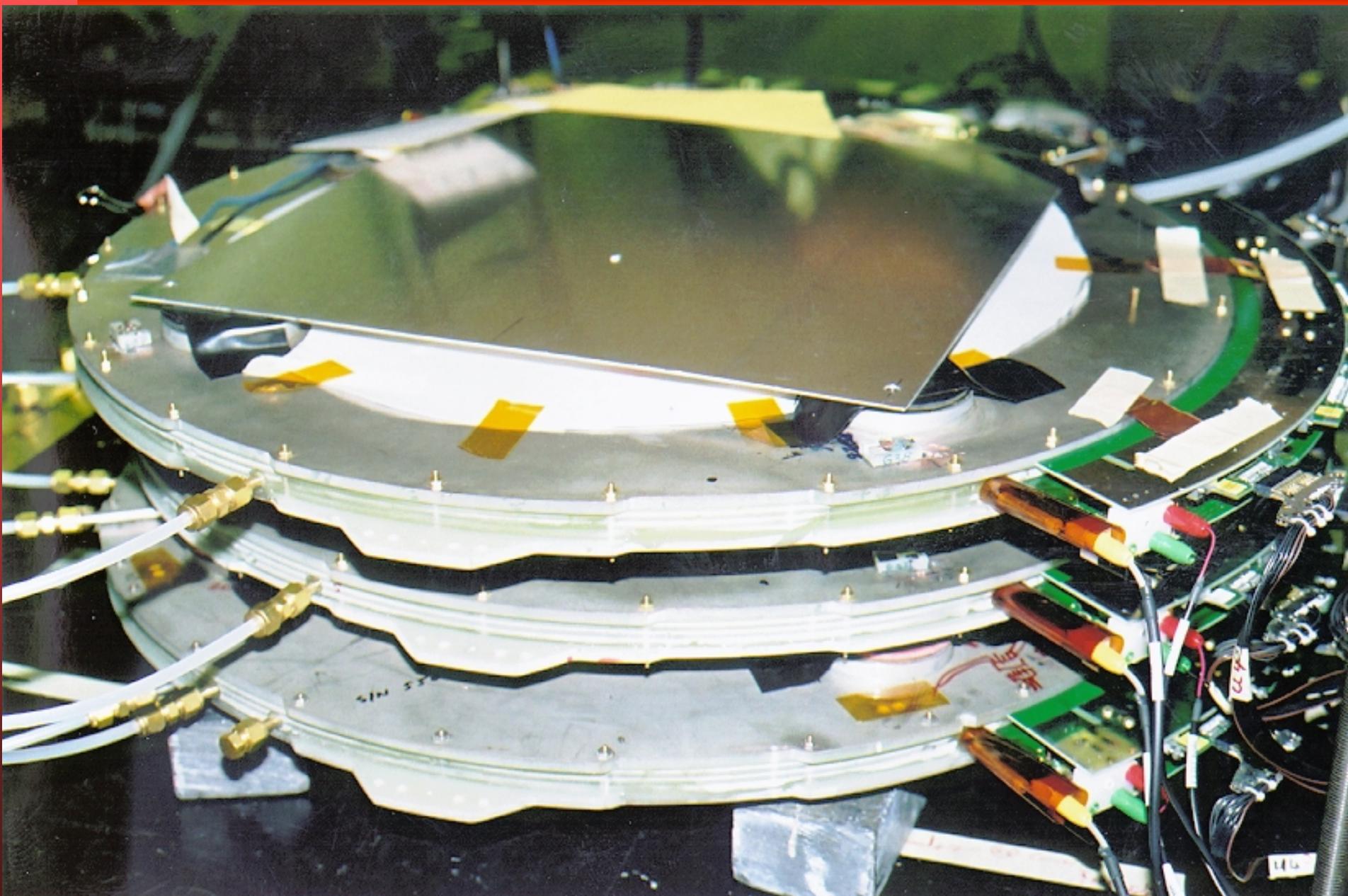
Magnet is cooling

Commissioning
begins this week

Mapping complete by
end of March

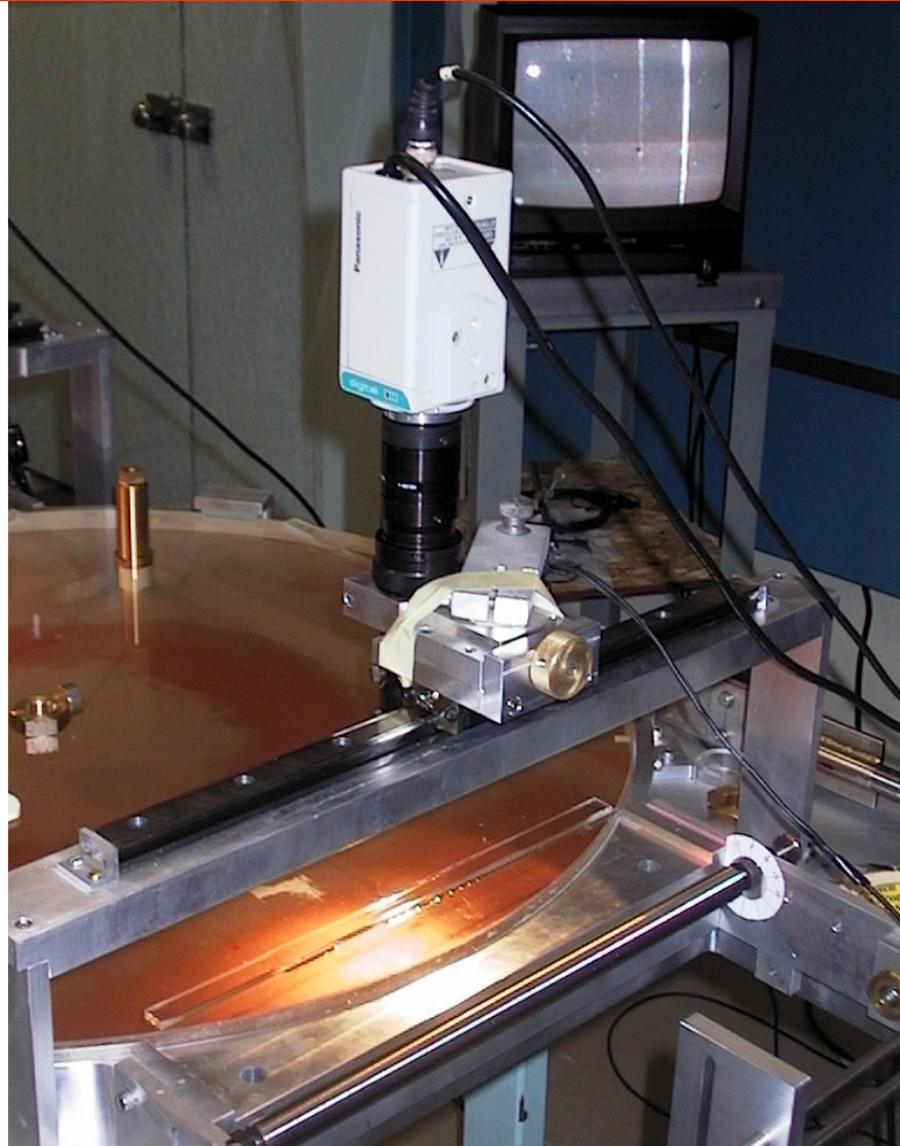


TWIST - Chambers

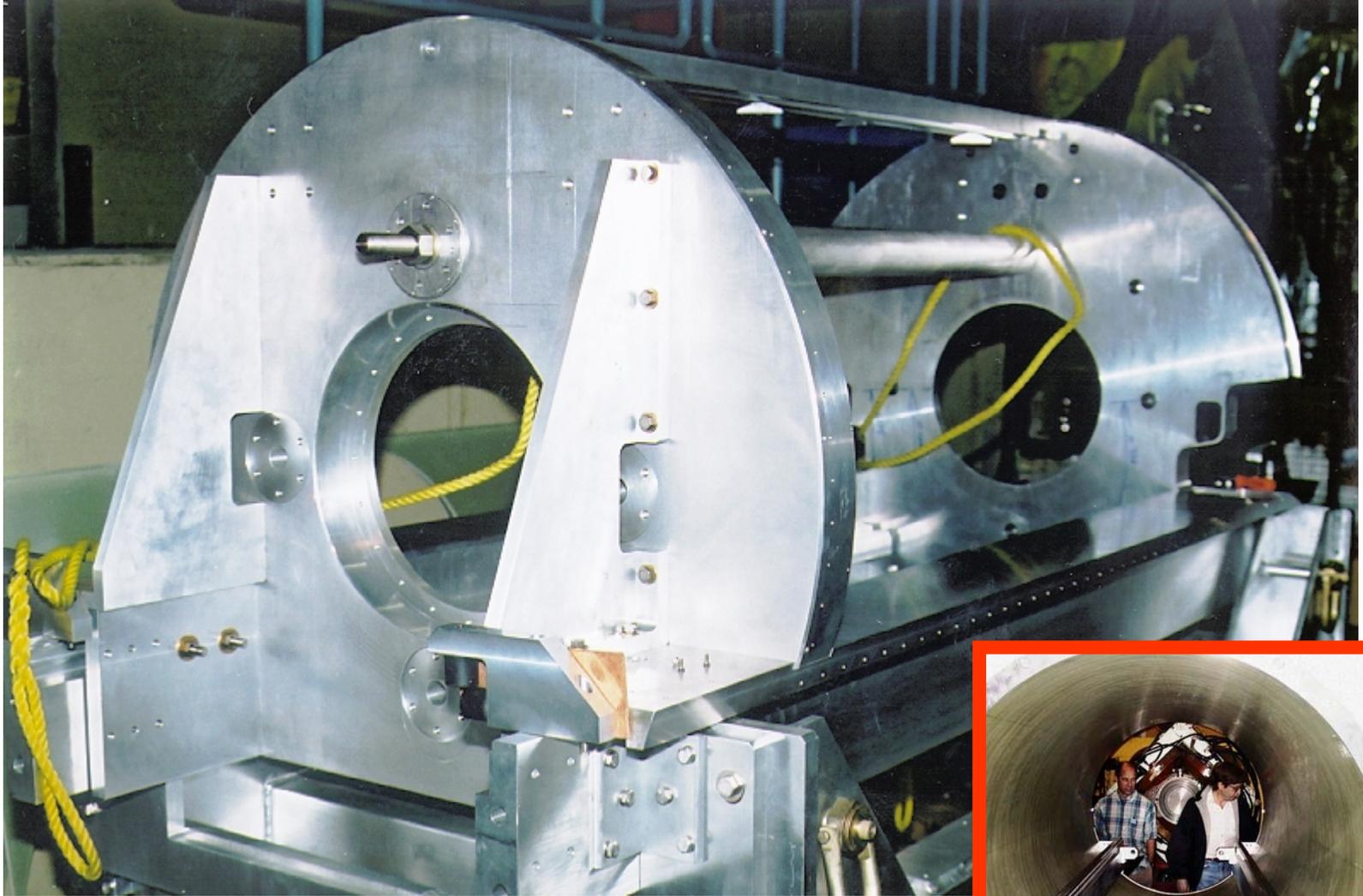


TWIST Glass Planes

Planes are assembled on glass plates with optical precision relative to longitudinal coordinate



TWIST – Chamber Support Cradle



Accepted Experimental Values

$$\rho = 0.7518 \pm 0.0026$$

$$P_{\mu\xi} = 1.0027 \pm 0.0085$$

$$\delta = 0.7486 \pm 0.0038$$

$$\eta = -0.007 \pm 0.013$$

E614 Proposal

$$\sigma_{\rho} = \pm 0.00005 \pm 0.00009$$

$$\sigma_{P_{\mu\xi}} = \pm 0.00010 \pm 0.00010$$

$$\sigma_{\delta} = \pm 0.00008 \pm 0.00010$$

$$\sigma_{\eta} \approx \pm 0.003$$

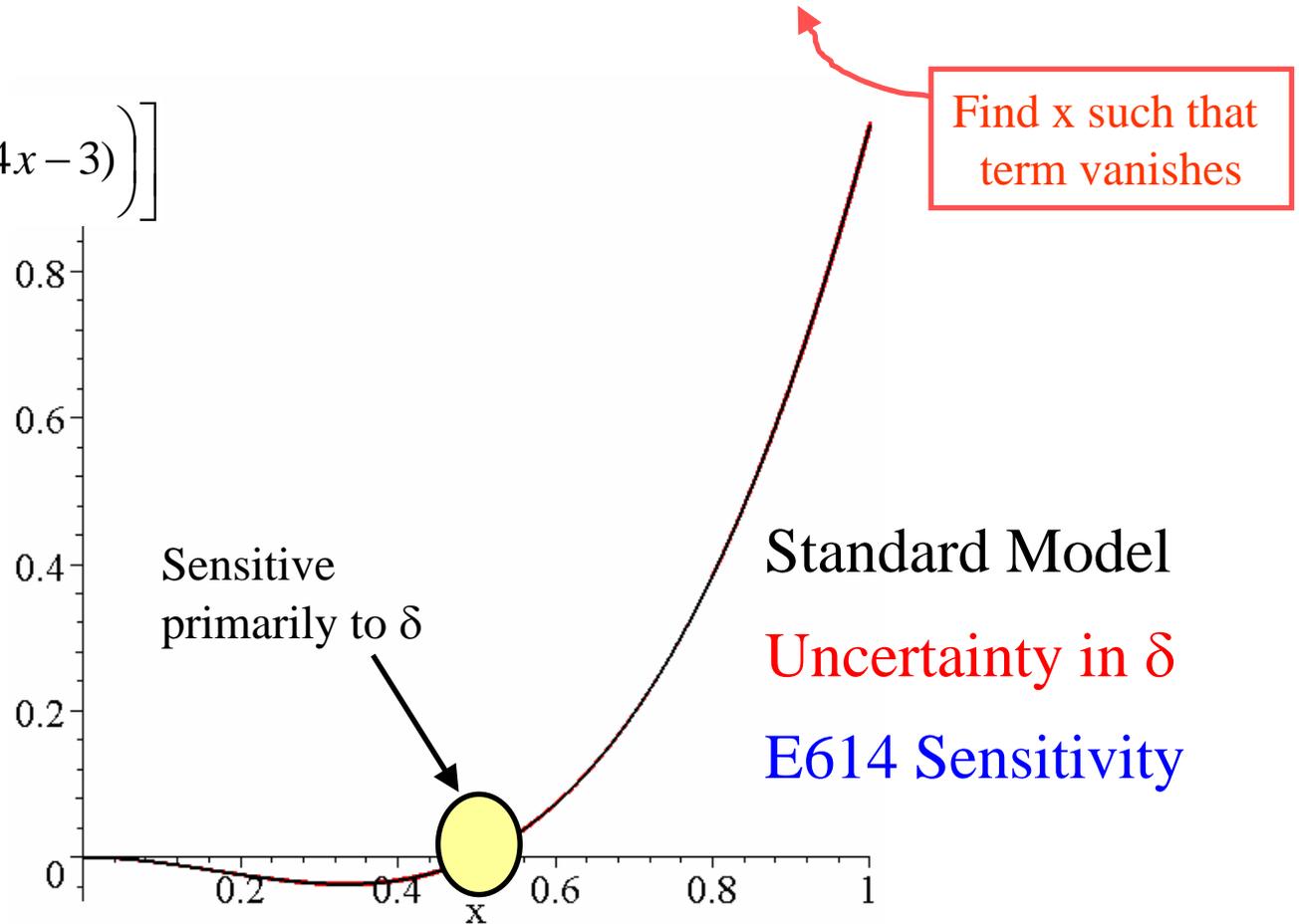
25-60 fold improvement in precision on the Michel parameters

3-10 fold improvement in couplings

The (forward - backward) distribution goes flat at a value of x dependant (only) upon δ

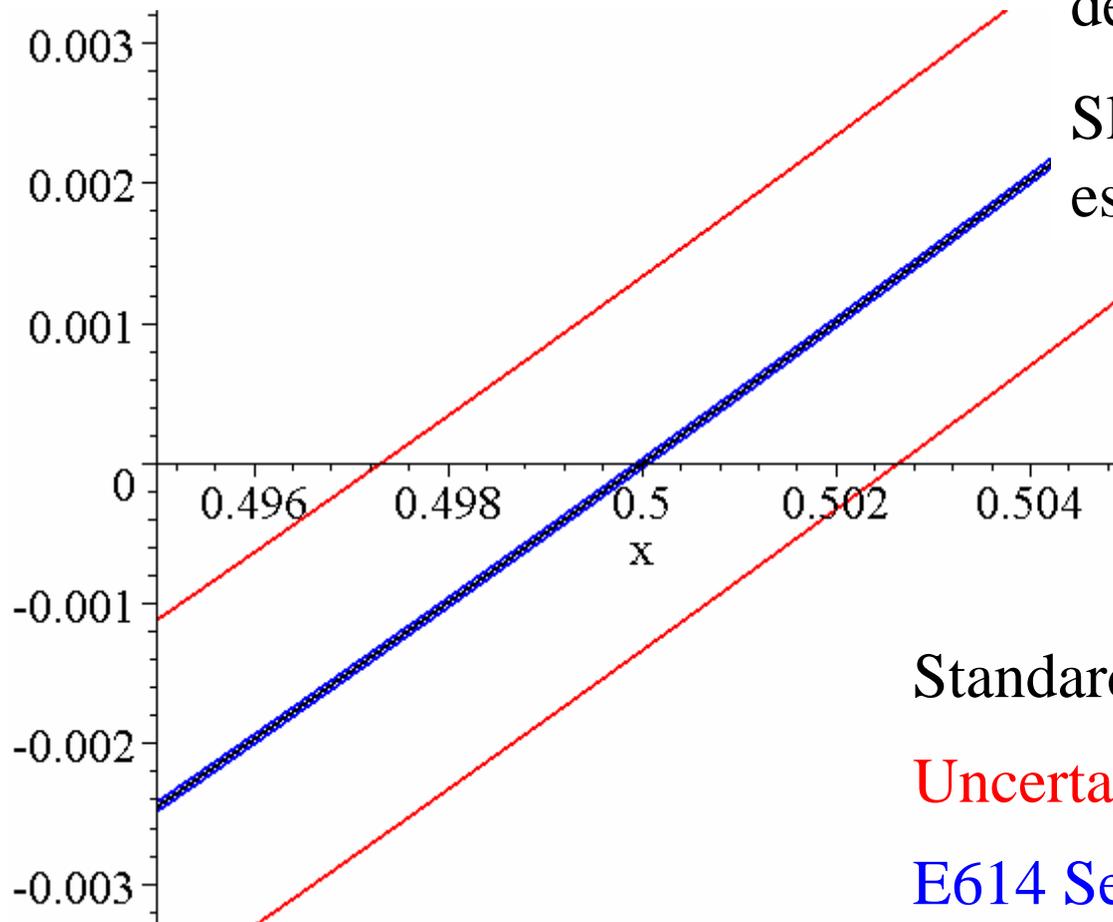
$$[Forward - Backward] \sim x^2 \left[2P_{\mu\xi} \cos(\theta) \left(1 - x + \frac{2}{3} \delta(4x - 3) \right) \right]$$

$$x^2 \left[2P_{\mu\xi} \left(1 - x + \frac{2}{3} \delta(4x - 3) \right) \right]$$



Same as the previous slide - on expanded scale

$$x^2 \left[2P_{\mu\xi} \left(1 - x + \frac{2}{3} \delta(4x - 3) \right) \right]$$



Zero crossing
determines δ

Slope is
essentially $P_{\mu\xi}$

Standard Model

Uncertainty in δ

E614 Sensitivity

Minimal extensions to the Standard Model

Allowing only vector couplings result in simplified Michel parameters

$$\rho \equiv \frac{3}{4} \left[|g_{LL}^V|^2 + |g_{RR}^V|^2 + |g_{LR}^T|^2 + |g_{RL}^T|^2 \right] \\ + \frac{3}{16} \left[|g_{LL}^S|^2 + |g_{RR}^S|^2 + |g_{LR}^S|^2 + |g_{RL}^S|^2 \right] \\ - \frac{3}{4} \left[\text{Re}(g_{LR}^S g_{LR}^{T*}) + \text{Re}(g_{RL}^S g_{RL}^{T*}) \right]$$

$$\xi \delta \equiv \frac{3}{4} \left[|g_{LL}^V|^2 - |g_{RR}^V|^2 - |g_{LR}^T|^2 + |g_{RL}^T|^2 \right] \\ + \frac{3}{16} \left[|g_{LL}^S|^2 - |g_{RR}^S|^2 - |g_{LR}^S|^2 + |g_{RL}^S|^2 \right] \\ - \frac{3}{4} \left[\text{Re}(g_{LR}^S g_{LR}^{T*}) - \text{Re}(g_{RL}^S g_{RL}^{T*}) \right]$$

In the context of the model,
Four parameters and four unknowns

$$\xi \equiv |g_{LL}^V|^2 + 3|g_{LR}^V|^2 - 3|g_{RL}^V|^2 - |g_{RR}^V|^2 + 5|g_{LR}^T|^2 \\ - 5|g_{RL}^T|^2 + \frac{1}{4}|g_{LL}^S|^2 - \frac{1}{4}|g_{LR}^S|^2 + \frac{1}{4}|g_{RL}^S|^2 - \frac{1}{4}|g_{RR}^S|^2 \\ + 4\text{Re}(g_{LR}^S g_{LR}^{T*}) - 4\text{Re}(g_{RL}^S g_{RL}^{T*})$$

$$\eta \equiv \frac{1}{2} \text{Re} \left[g_{LL}^V g_{RR}^{S*} + g_{RR}^V g_{LL}^{S*} \right] \\ + \frac{1}{2} \text{Re} \left[g_{RL}^V (g_{LR}^{S*} + 6g_{LR}^{T*}) + g_{LR}^V (g_{RL}^{S*} + 6g_{RL}^{T*}) \right]$$

Anticipated sensitivity to new couplings

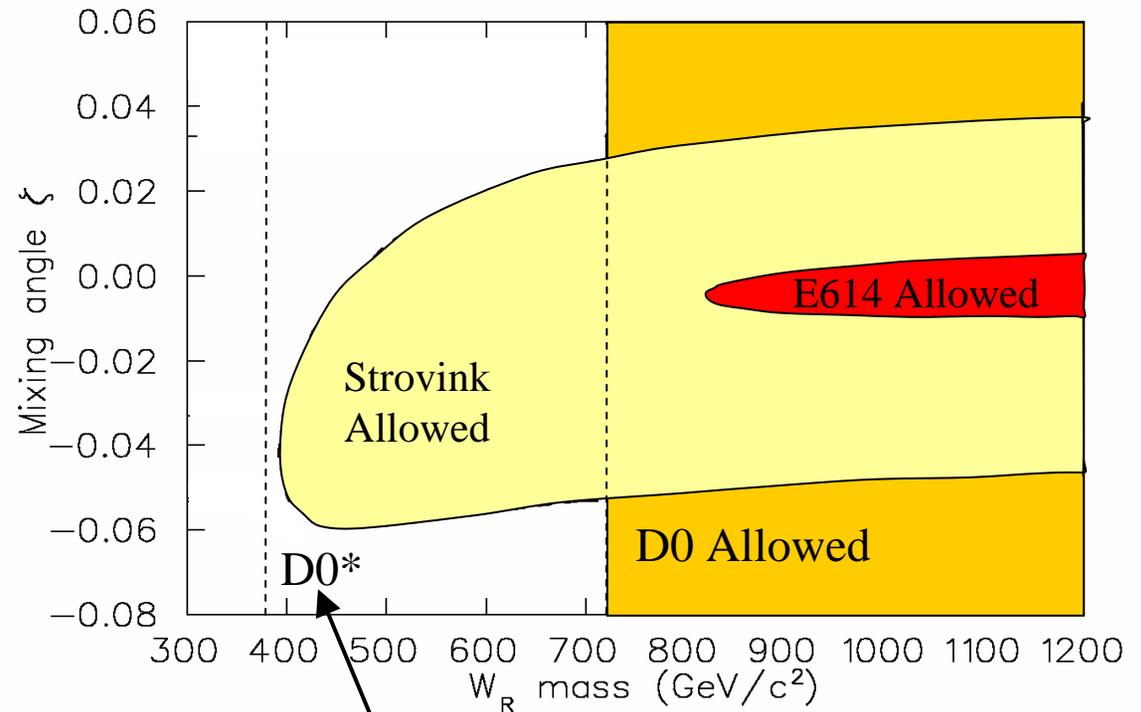
	Current Limits	E614(A)	E614(B)	E614(C)	E614(D)
$ g_{RR}^S $	<0.066	—	—	0.020	0.045
$ g_{RR}^V $	<0.033	0.012	0.014	0.013	0.022
$ g_{LR}^S $	<0.125	—	—	0.027	0.046
$ g_{LR}^V $	<0.060	0.012	0.013	0.012	0.018
$ g_{LR}^T $	<0.036	—	0.009	—	0.013
$ g_{RL}^S $	<0.424	—	—	—	—
$ g_{RL}^V $	<0.110	0.012	0.012	0.011	—
$ g_{RL}^T $	<0.122	—	0.008	—	—
$ g_{LL}^S $	<0.55	—	—	—	—
$ g_{LL}^V $	>0.96	>0.99977	>0.99953	—	—

Upper limits (90% CL) for weak coupling constants with current limits taken from the Particle Data Group. Improved limits expected from TWIST based on measurements of ρ , ξ , δ and η assume:

- (A) V, A couplings only,
- (B) V, A and T couplings,
- (C) V, A and S couplings or
- (D) most general V, A, S, and T derivative-free couplings.

One way of looking at the discovery potential

Assume manifest L-R
Symmetry
ie $g_R = g_L$
 $CKM_R = CKM_L$
and no cp violation



$$V_{ud}^R = 0$$

$$V_{us}^R = 1$$

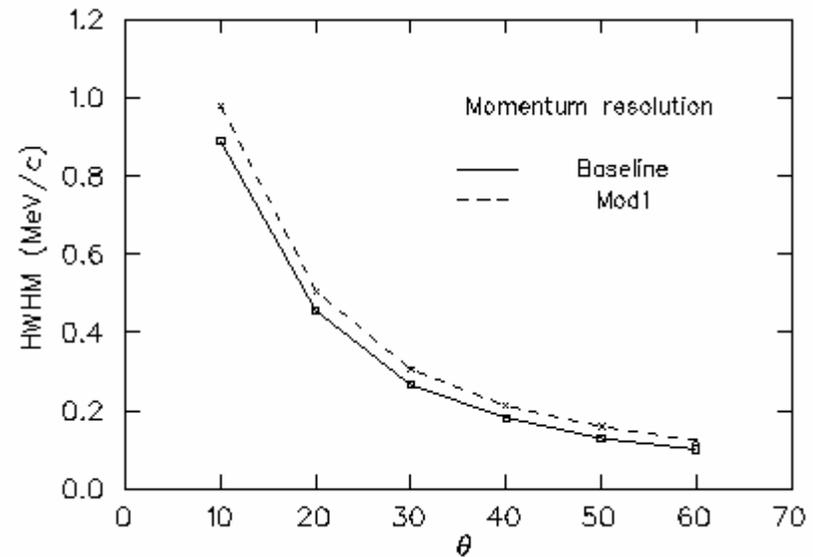
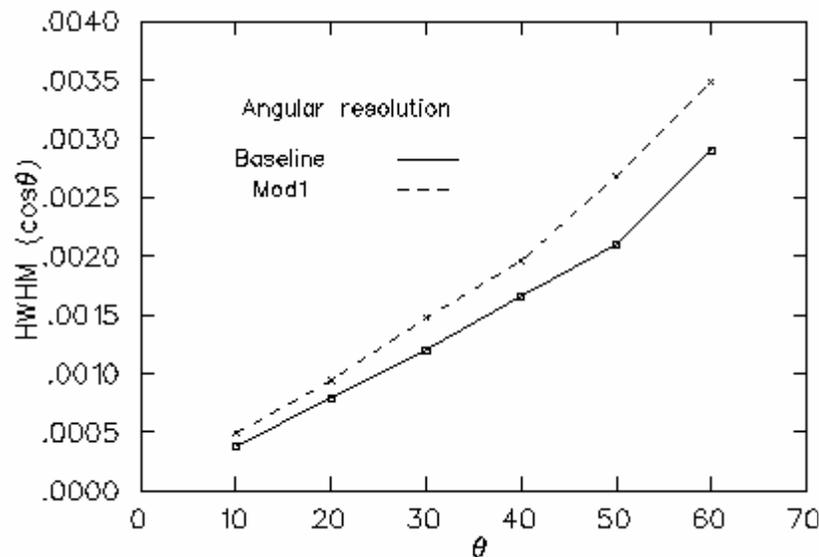
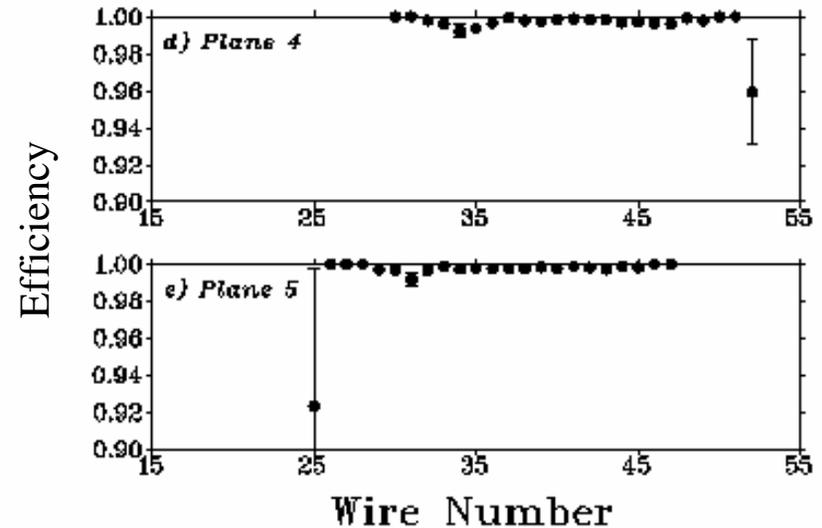
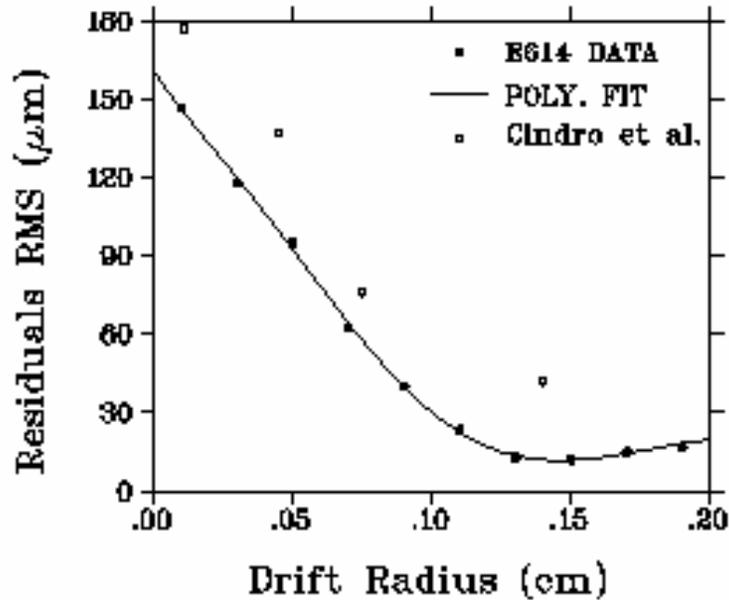
Beta decay, $p\bar{p}$ direct production, and muon decay
are complimentary

E614 Timeline

- ✓ **High Priority at TRIUMF – 1993**
- ✓ **First Capital Funding – April 1997**
- ✓ **WC Review - January 1999**
- ✓ **Mechanical Review - June 1999**
- ✓ **Beam Tests - final prototype - August 1999**
- ✓ **Full WC Production underway - March 2000**
- **WC Module Completion May 2000 – April 2001**
- ✓ **WC Bench tests beginning June 2000**
- ✓ **Yoke assembly December 2000**
- **Yoke, Solenoid, and cryogenics Commissioning: February - April 2001**
- **First beam: Summer of 2001**

- **Preliminary Physics: December 2002**

Spectrometer Resolution



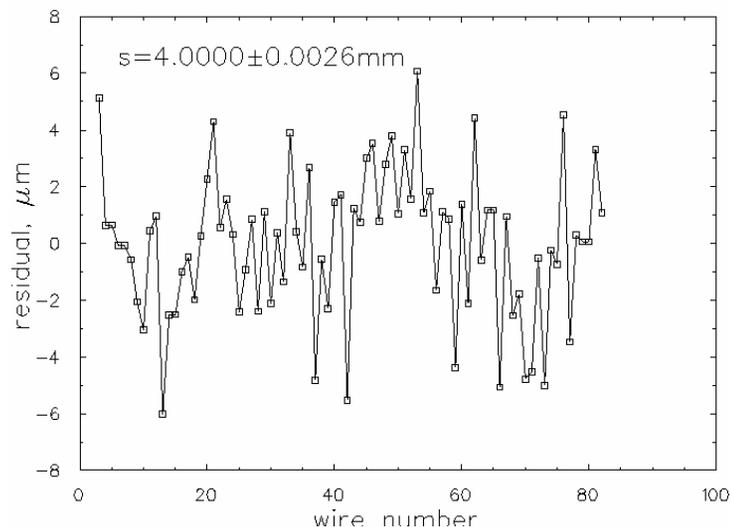
TWIST - Chambers

Quality Control on stringing of Wire Planes

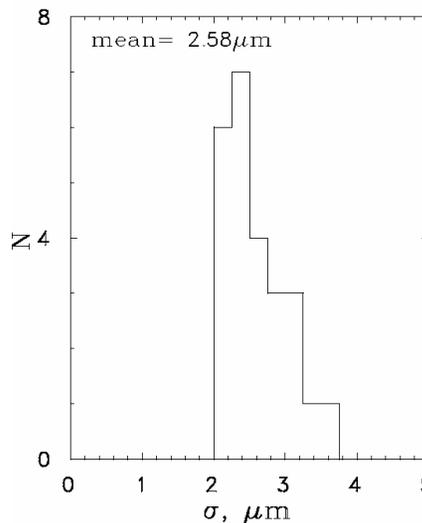
The figures show:

1. Wire-to-wire variation in z position for a typical plane;
 $\sigma = 2.6 \mu$
2. Average error in wire position over 25 drift planes;
 $\sigma = 2.58 \mu$
3. Average wire tension over 38 drift planes;
rms = 0.94g

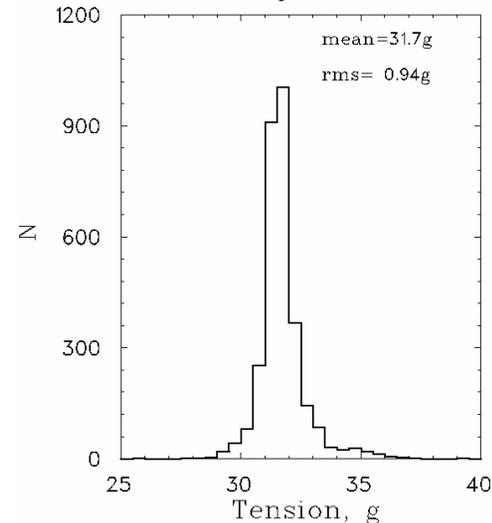
Plane 035DC, sense wires



Precision of DCs wire positioning
25 planes, centers of chambers



Tension of sense wires
38 DC planes



TWIST - Electronics

TWIST Requires

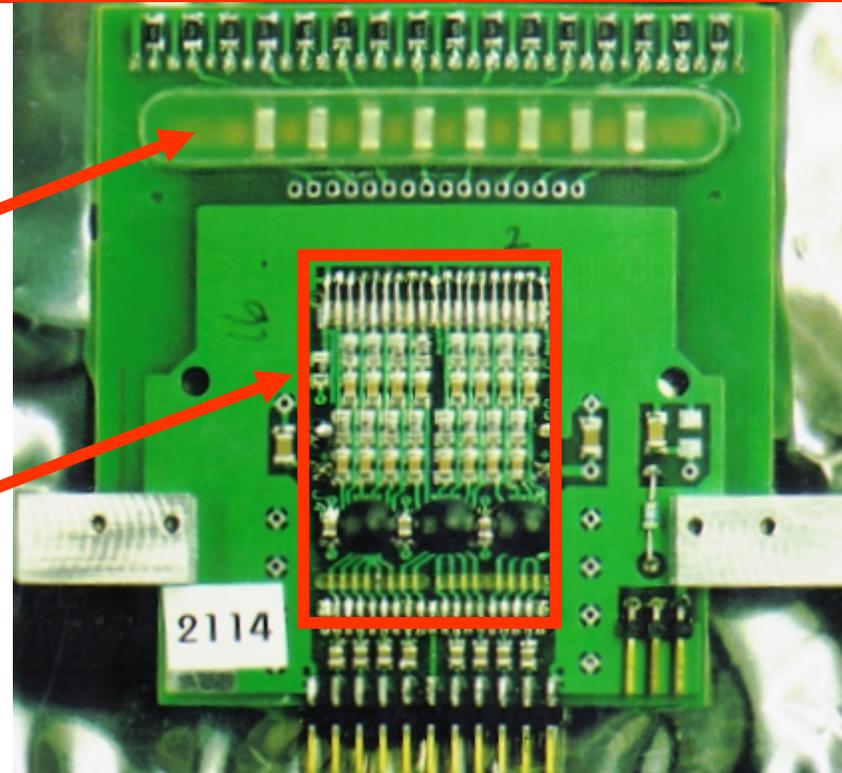
- 240 preamplifiers
- 268 postamplifiers
- 42 TDC's

Status

- 86 preamplifiers tested, 41 in mid-production
- 120 postamplifiers tested, 180 more in production
- 47 TDC's in hand

HV potting
and isolation
capacitors

VTX board



TWIST preamplifier

16 and 24 channel versions based on Fermilab CDF VTX boards

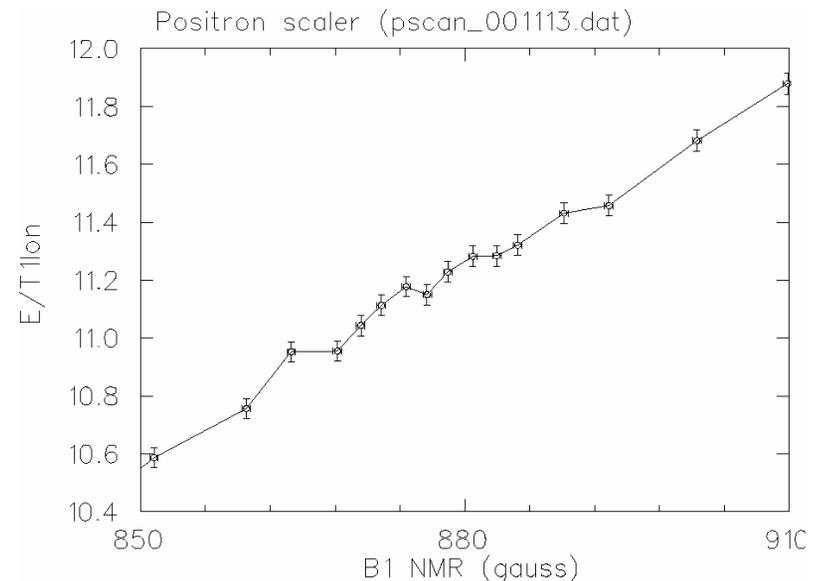
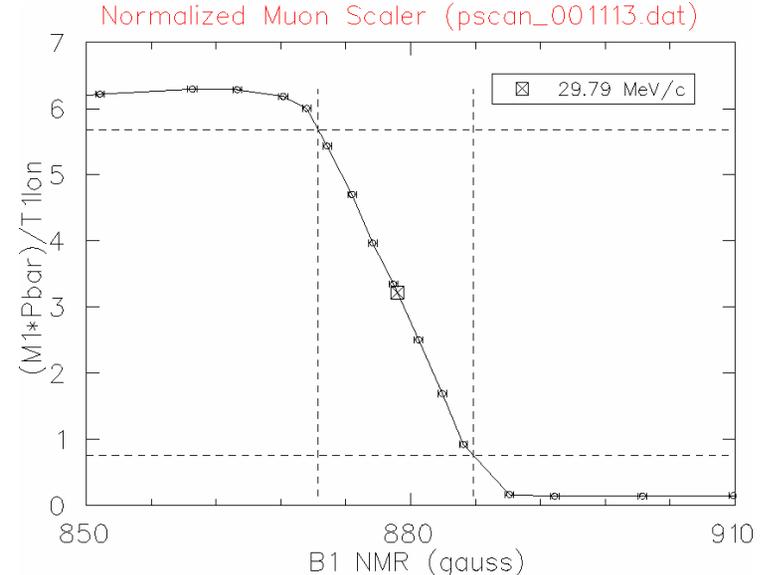
Cross talk is minimal (0.8% amplitude), and is easily rejected in software by cutting on pulse width

TWIST – Positron Background

Beamline studies from
October/November 2000

- Backgrounds
 - Rates: $e^+/\mu^+ \sim 4$
(as expected)

A pyrolytic graphite target will
give us a 33% improvement in
the rate relative to the
positrons



TWIST – RF Cuts

Flight time through beamline

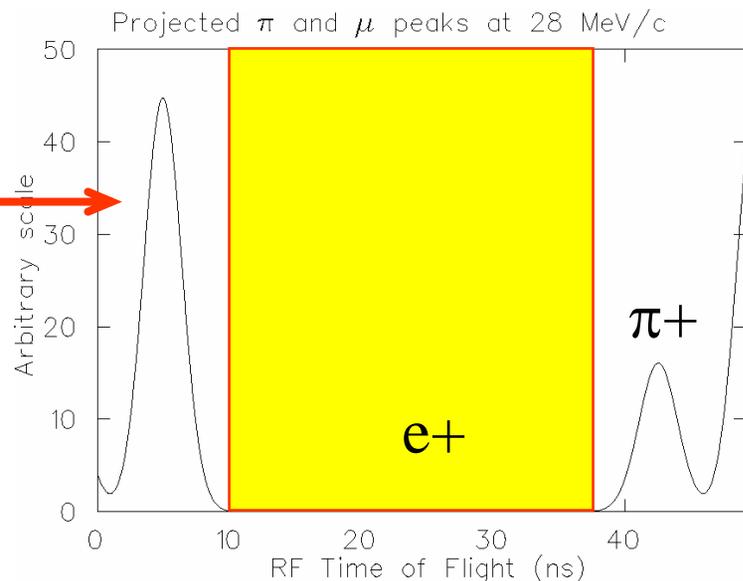
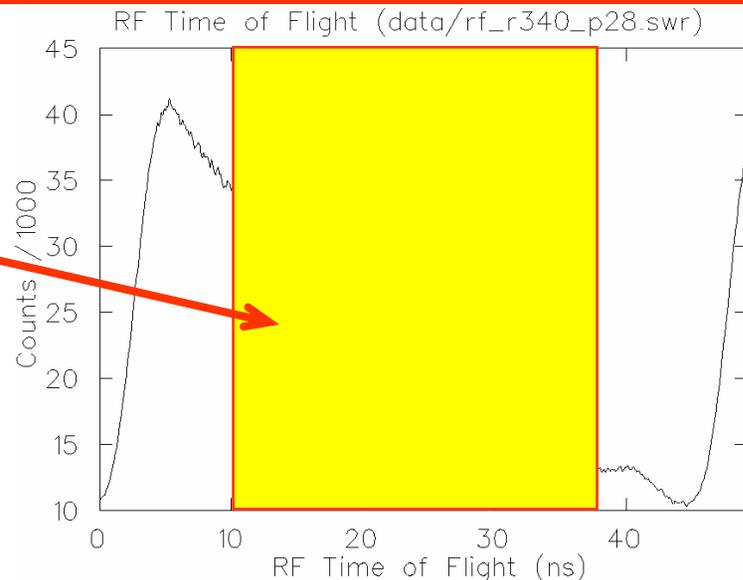
Surface Muons gated on
cyclotron RF

Time characteristic of π decay

Backgrounds (extrapolated
from higher momentum)

Cloud Muons

Rate: 9% that of
surface muons

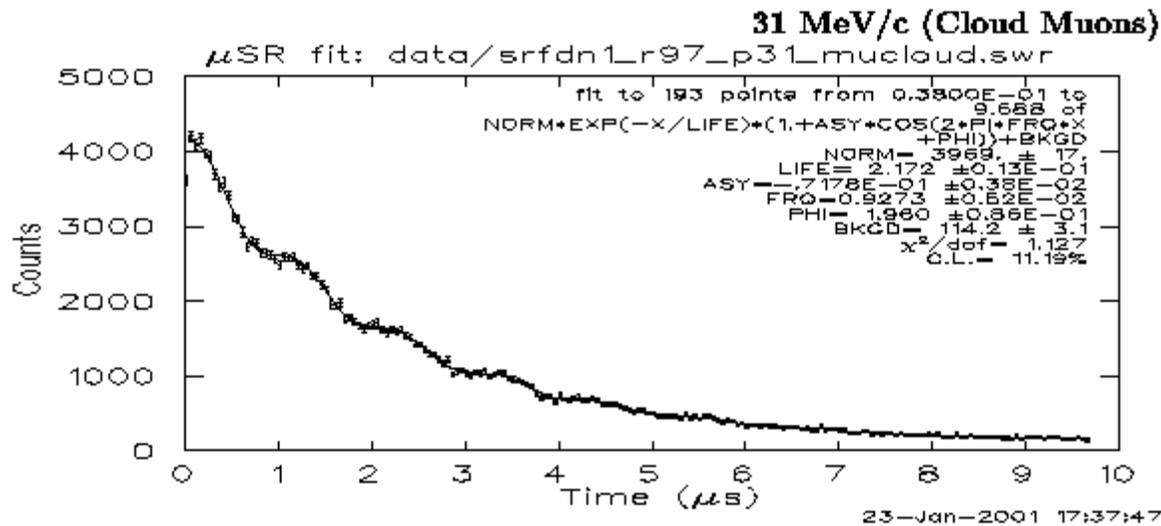
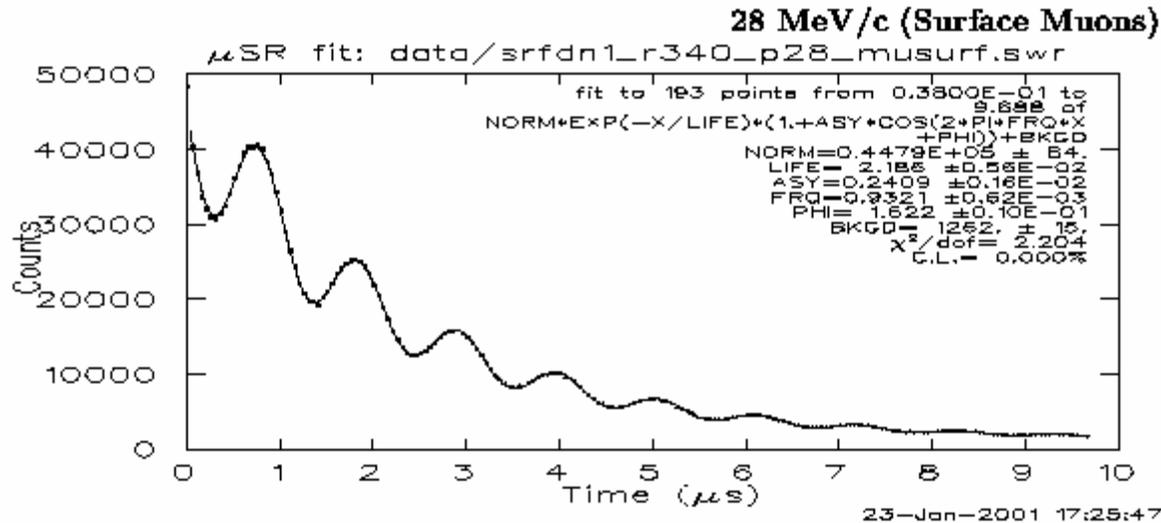


TWIST – Cloud muon polarization

Surface muons

Polarization of the cloud muons is approximately 0.30 (opposite to the surface muon polarization of -1.0)

Cloud muon flux is 9% that of the surface muons

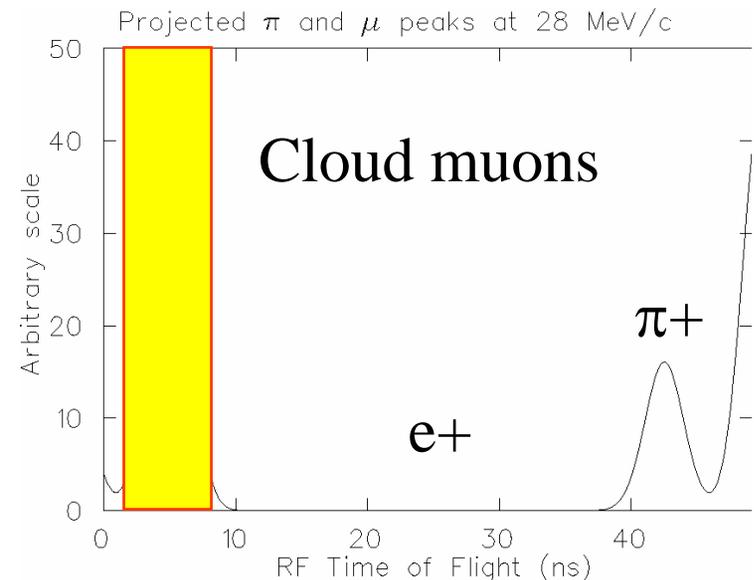
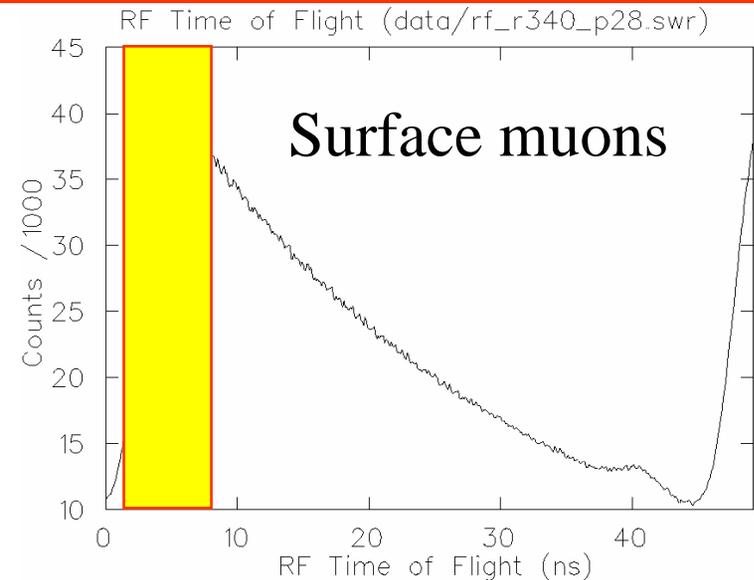


(Rob MacDonald – MSc data)

TWIST – Unpolarized data sample

Flight time through beamline

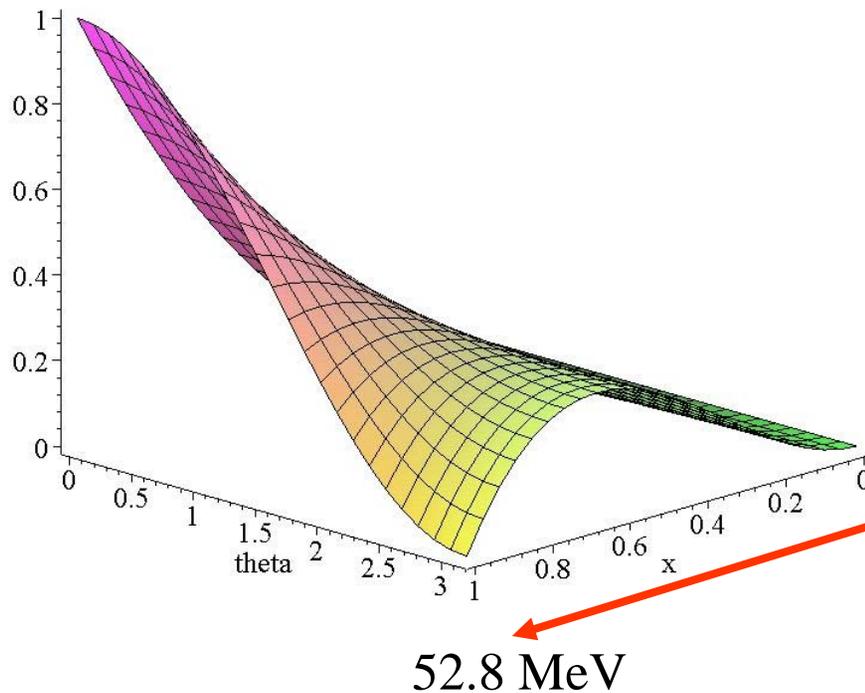
By selecting a data sample with an appropriate RF gate, we can select an unpolarized sample of muons



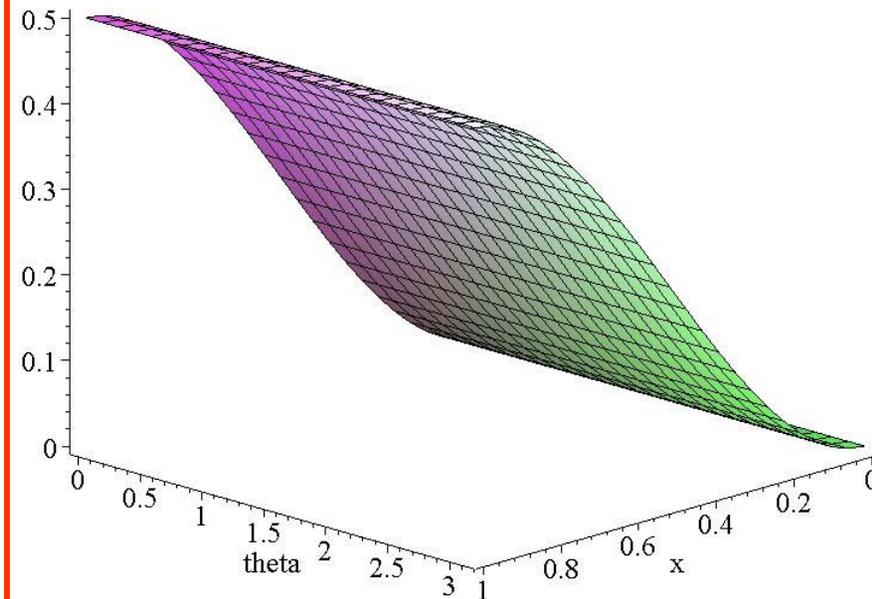
TWIST – Energy Calibration

The edge of the distribution is used to calibrate the energy scale at large x

The polarized distribution has no edge at forward angles



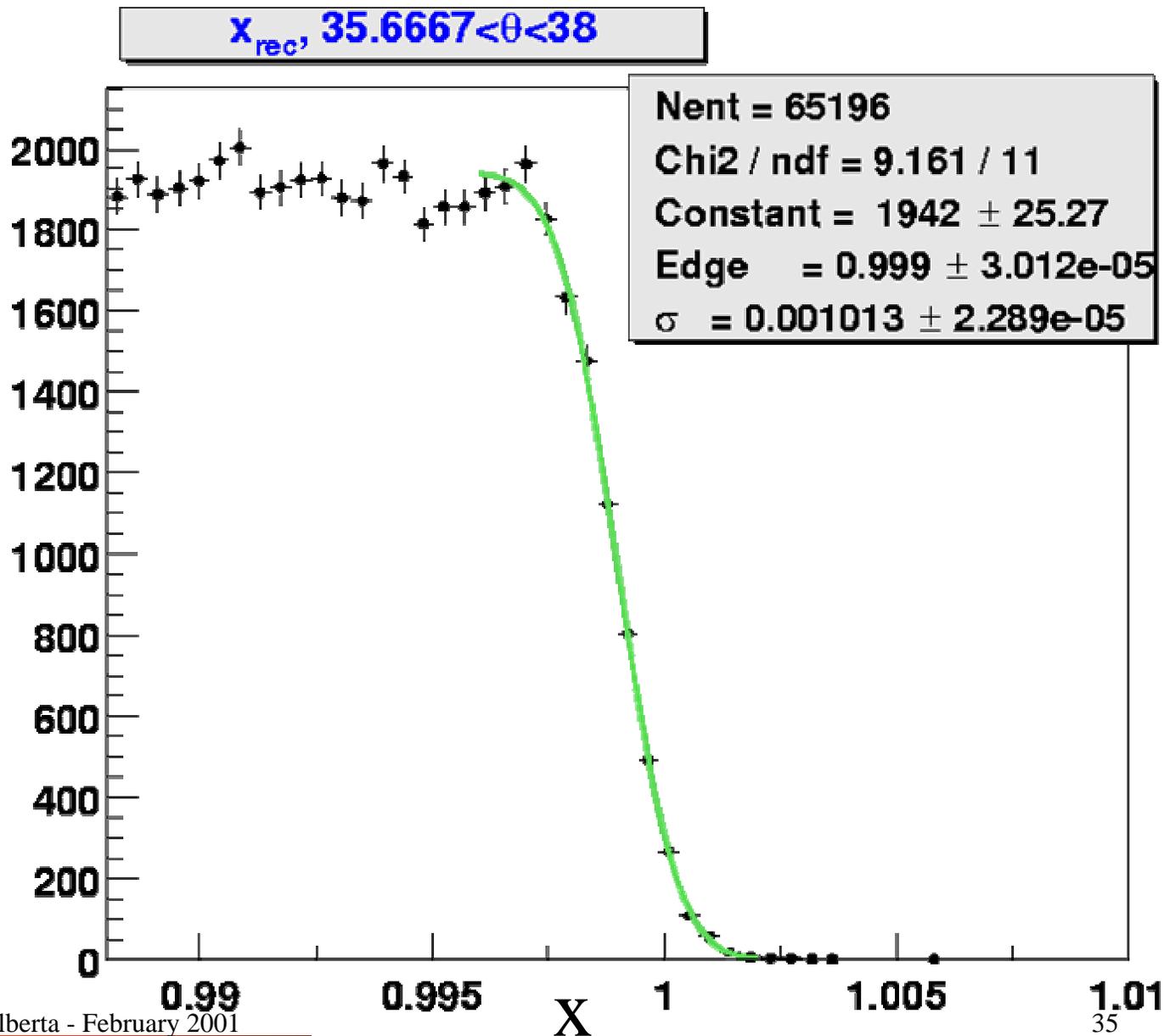
The unpolarized distribution will be used to calibrate the energy scale at all angles



TWIST – Energy Calibration

Endpoint energy calibrations can be done to a precision of approximately 2 keV (where ~10 keV is needed).

Unpolarized beam will be used to provide energy calibrations independent of angle.

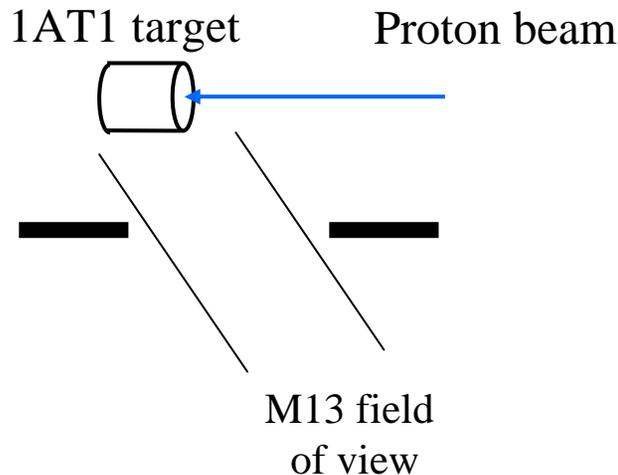


TWIST-1AT1 modifications

The surface muon beam is produced in part on the surface at which the protons enter, and in part along the length of the target cylinder.

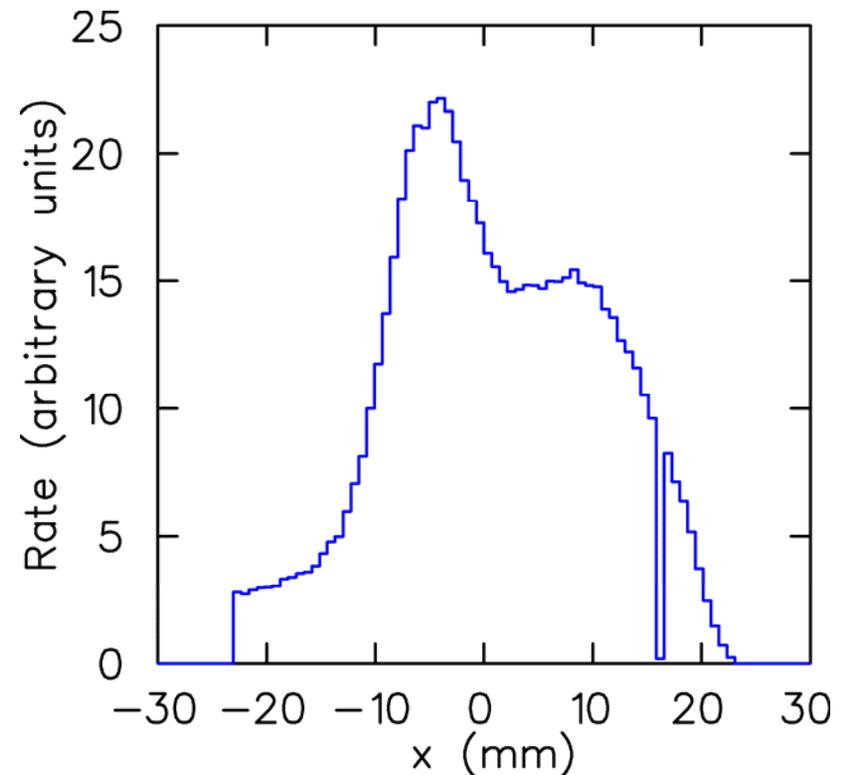
A shorter target would reduce the size of the beam spot

A hidden proton entry point would reduce sensitivity to wander in the proton beam

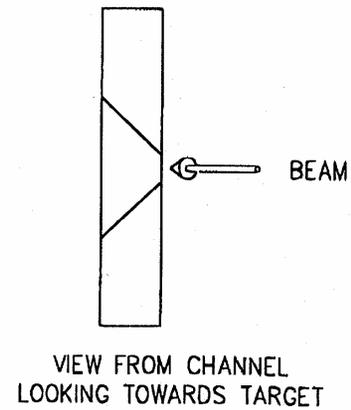
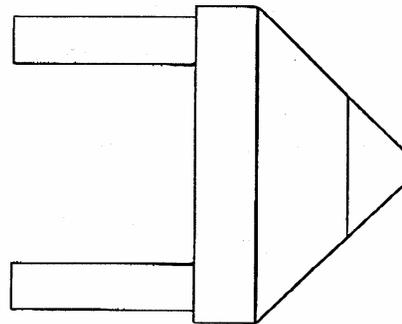
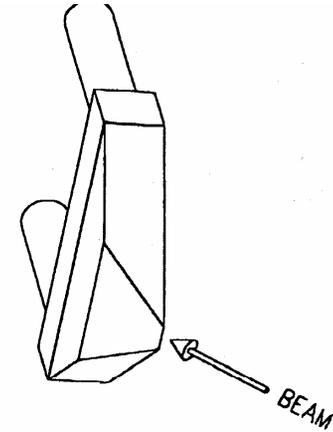
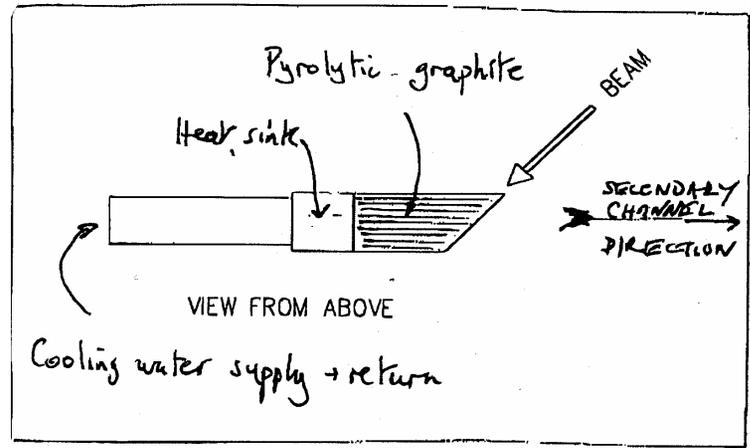


1AT1 target as imaged by M13

Surface muons



TWIST- Modified target



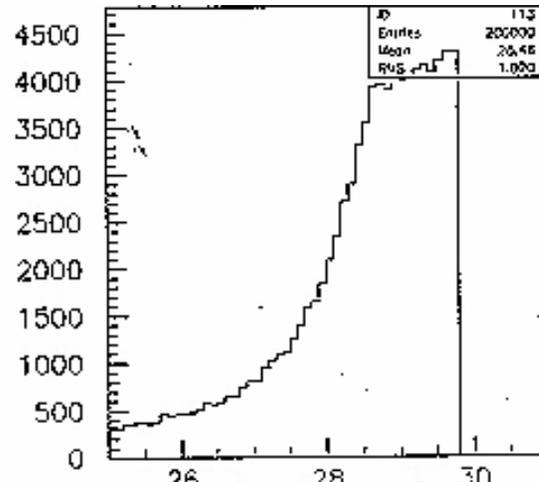
TWIST-Goals

- Summer 2001** - **Commissioning data. Preliminary alignments and calibrations**
- End of 2001** - **Michel distributions on tape suitable for preliminary determination of ρ and δ**
- 2002** - **Installation of the TEC**
 - **Modified production target**
 - **Beamline improvements, including realignments**
 - **Improved Michel distributions based upon experience with alignments and calibrations**
 - **Field alignment studies**
- 2003** - **Studies of depolarization in the stopping target**
 - **Preliminary $P_{\mu\xi}$ data**
 - **Precision measurements of ρ , δ and η**

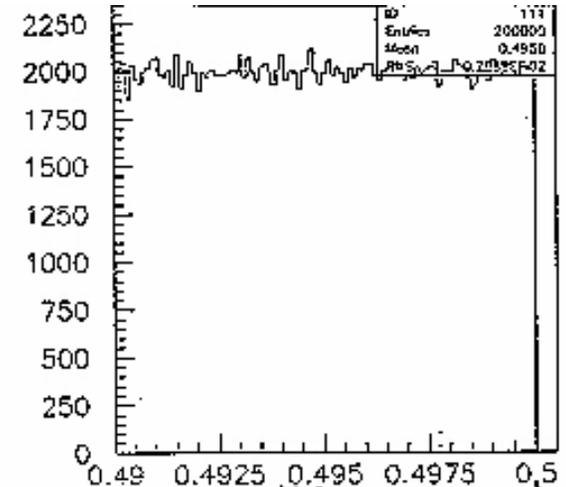
TWIST-1AT1 surface selection

$\Delta p/p$ of 1% selects muons from within about 20 microns of the surface.

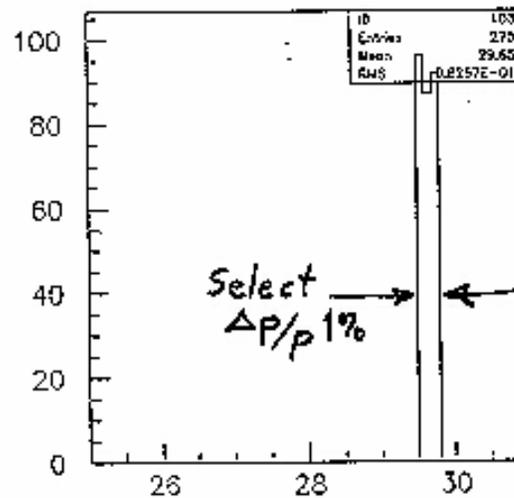
These muons have limited multiple scattering, and little depolarization



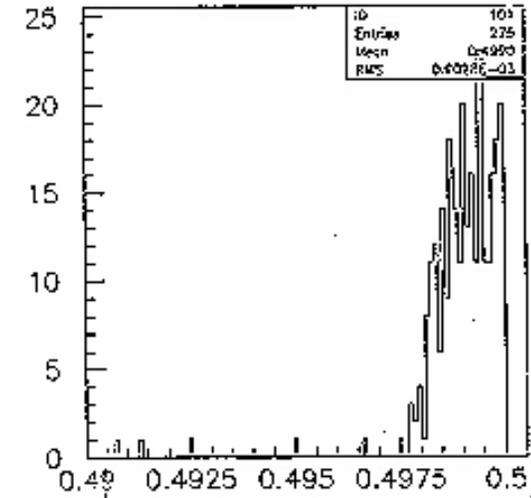
Momentum of μ^+ (Target Stop decay)



Stop(Target) pi decay positions



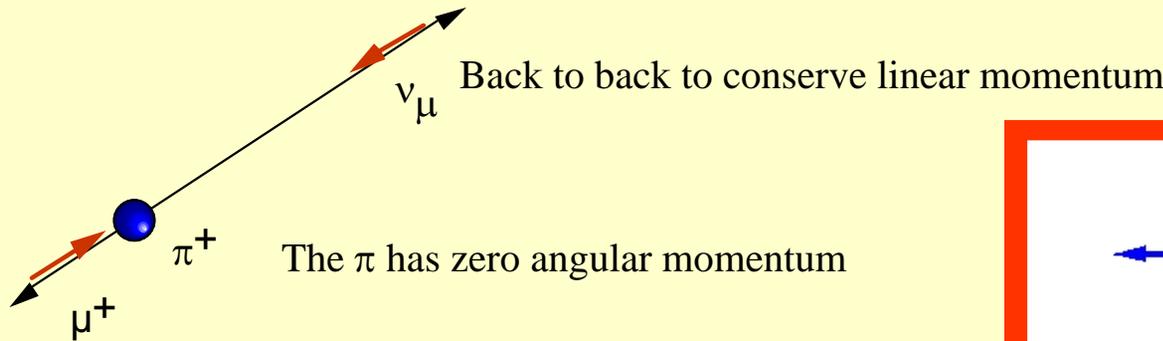
Momentum of μ^+ (Triggered Target Stop decay)



Stop(Target) pi decay positions

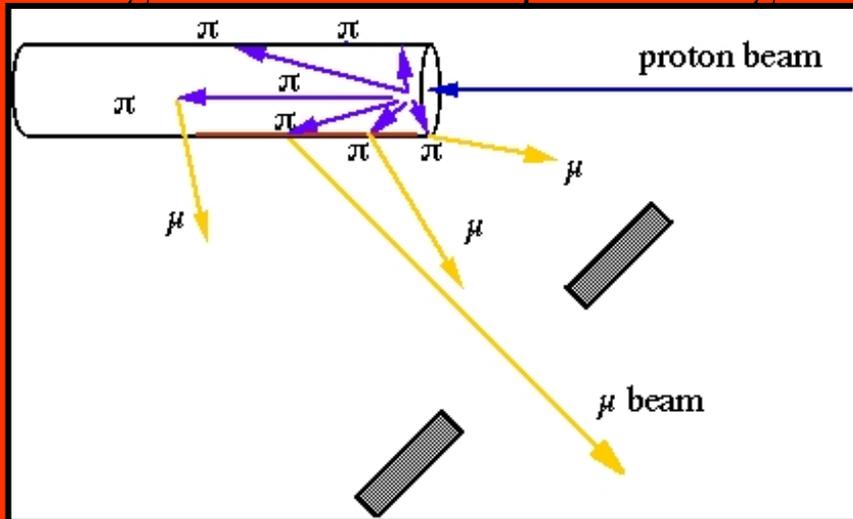
Secondary beams at TRIUMF

μ polarization due to 2-body decay

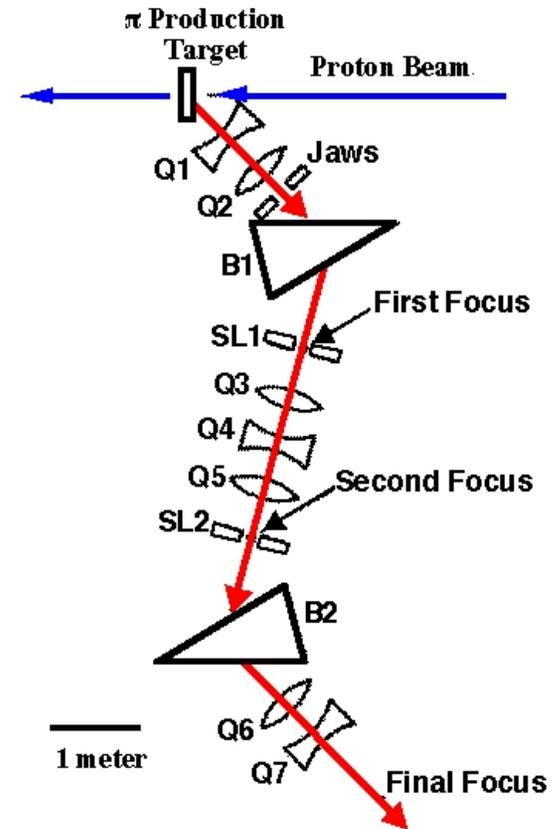


=> no angular momentum in the final system

μ selected from the surface of the production target suffer little multiple scattering



Channel resolution $\sim 1\%$ allows selection of μ produced within 25 microns of target surface



Momentum Resolution $\Delta p/p = 1\%$

TWIST – average energy loss

Planar chambers give us a simple dependence of energy loss on $1/\cos(\theta)$.

Each successive curve is the result of a track fit using only four successive chambers.

The difference between successive curves demonstrates the small incremental energy loss per plane of ~ 10 keV at 0 degrees

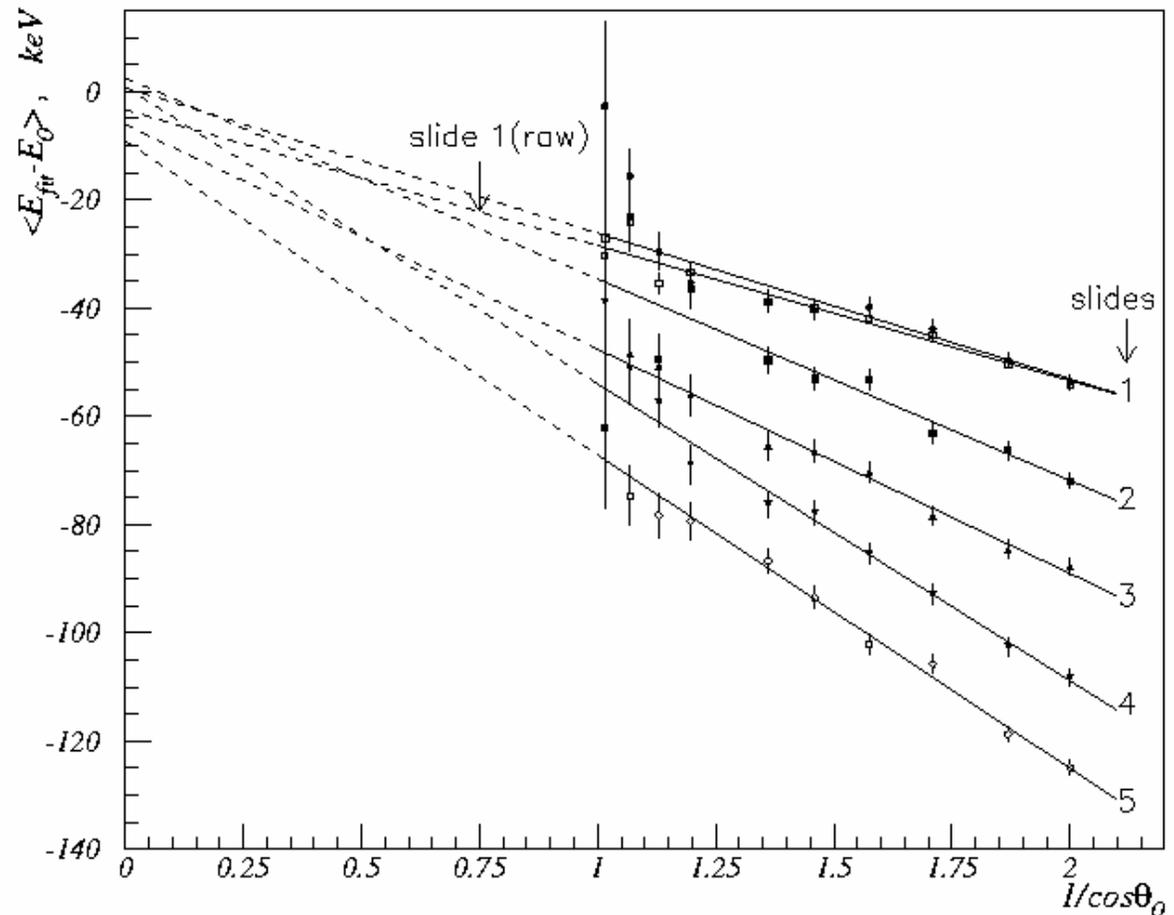
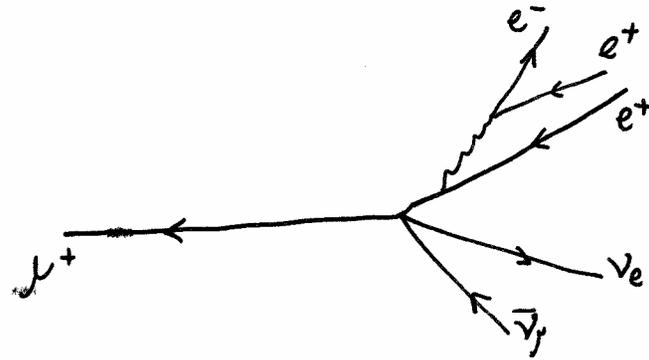
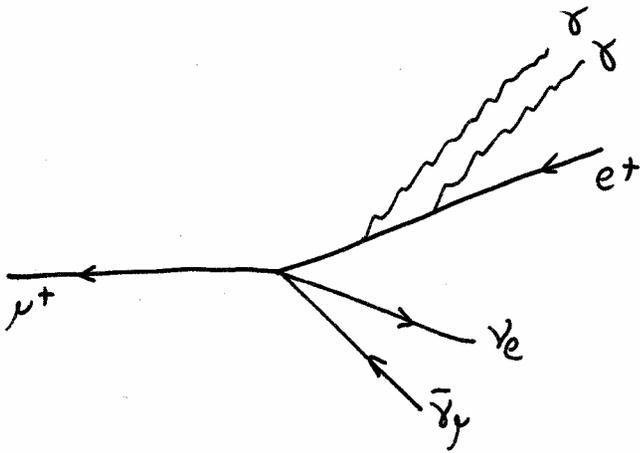


Fig. 9. Mean reconstructed positron energy \overline{E}_{fit} as a function of $1/\cos\theta$. $E_0=50$ MeV, $\sigma_{PDC}=50$ μm . Straight lines are the fits with function $\overline{E}_{fit} = E_0 - \alpha/\cos\theta$.

TWIST- Radiative corrections



Diagrams responsible for the double-logarithmic $O\left(\frac{\alpha^2}{\pi} \ln^2 \frac{m_\mu}{m_e}\right)$ corrections to the e^+ spectrum.

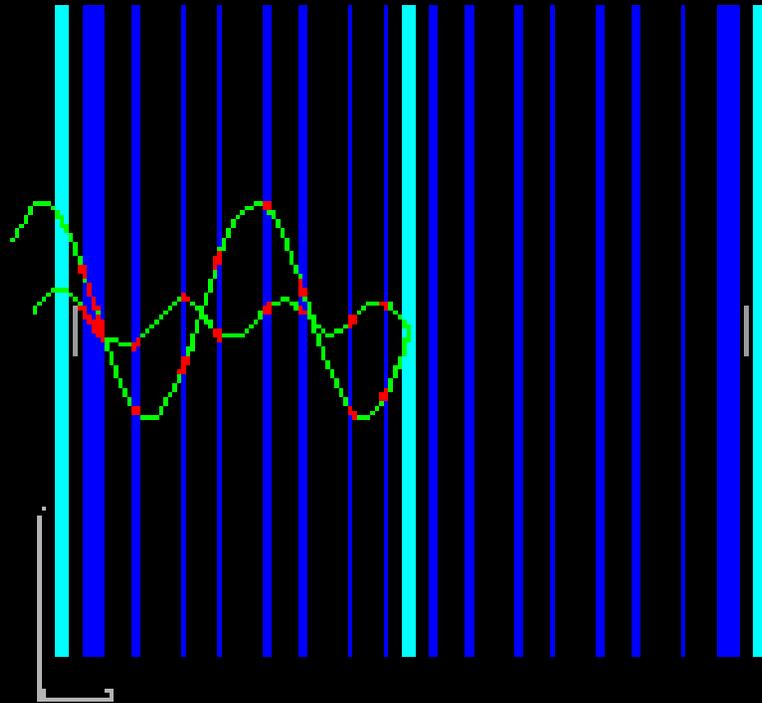
These diagrams have recently been calculated by Czarnecki and Arbusov (Alberta)

Happy physicist with magnet in hand

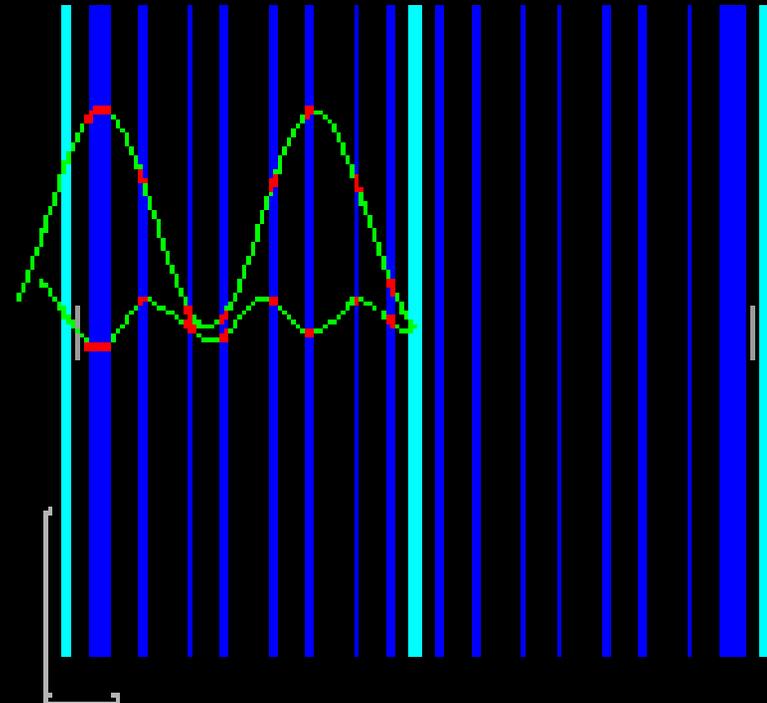


TWIST-Event Display

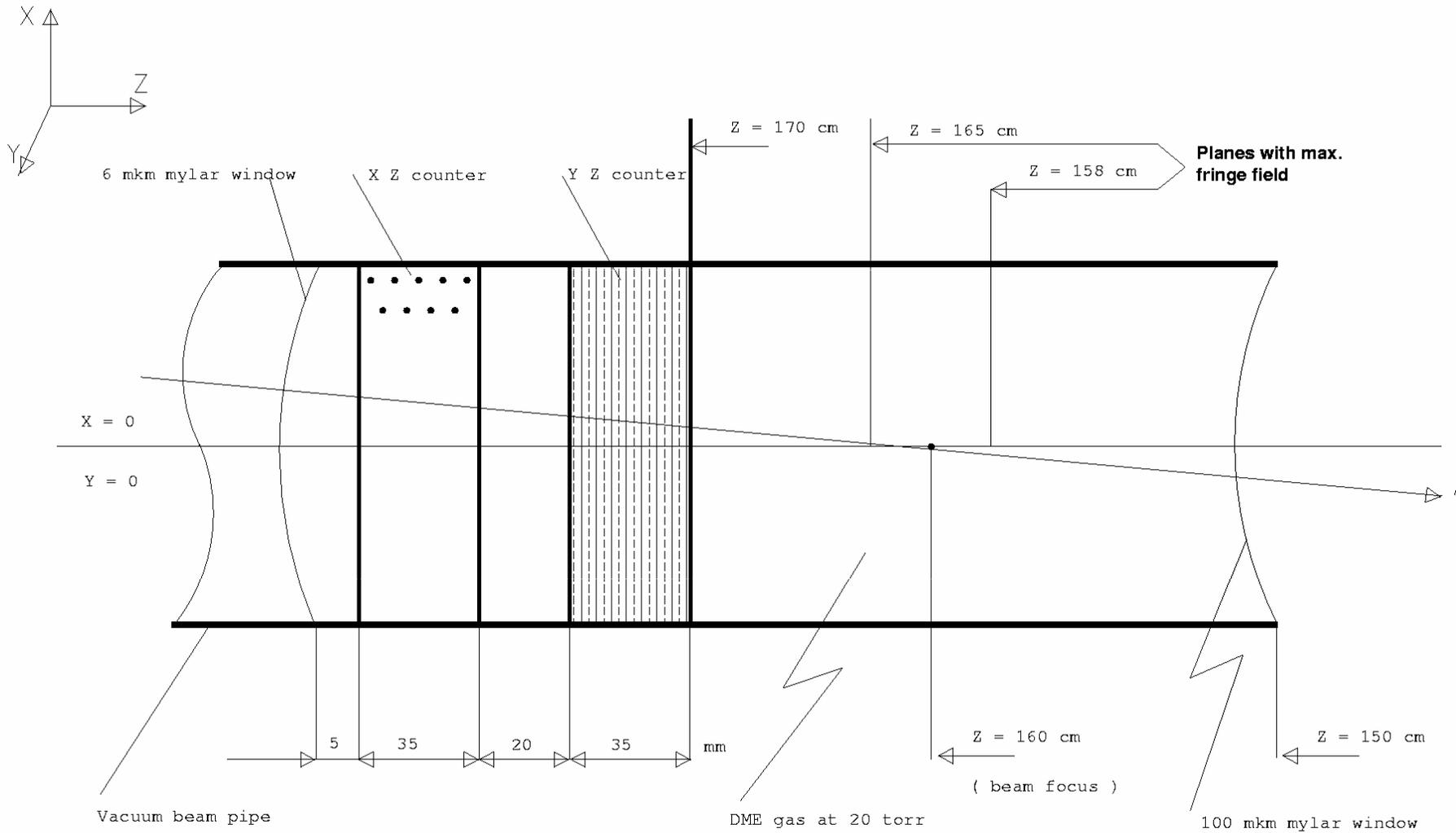
Top



Side



TWIST-TEC Design Concept



TWIST-TEC Design Concept

The TEC has been part of our planning since June 1998

Installation planned for Spring 2002

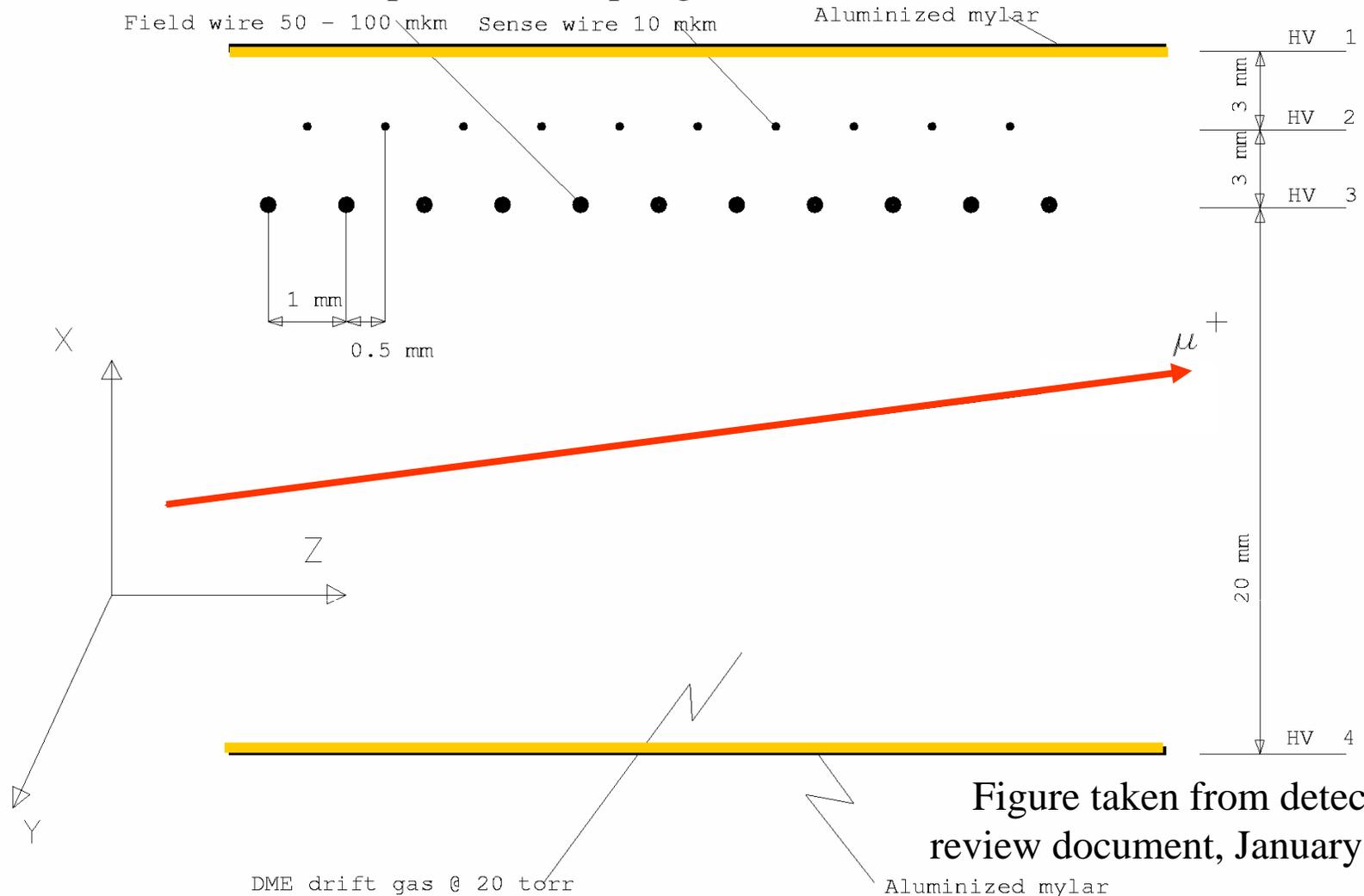
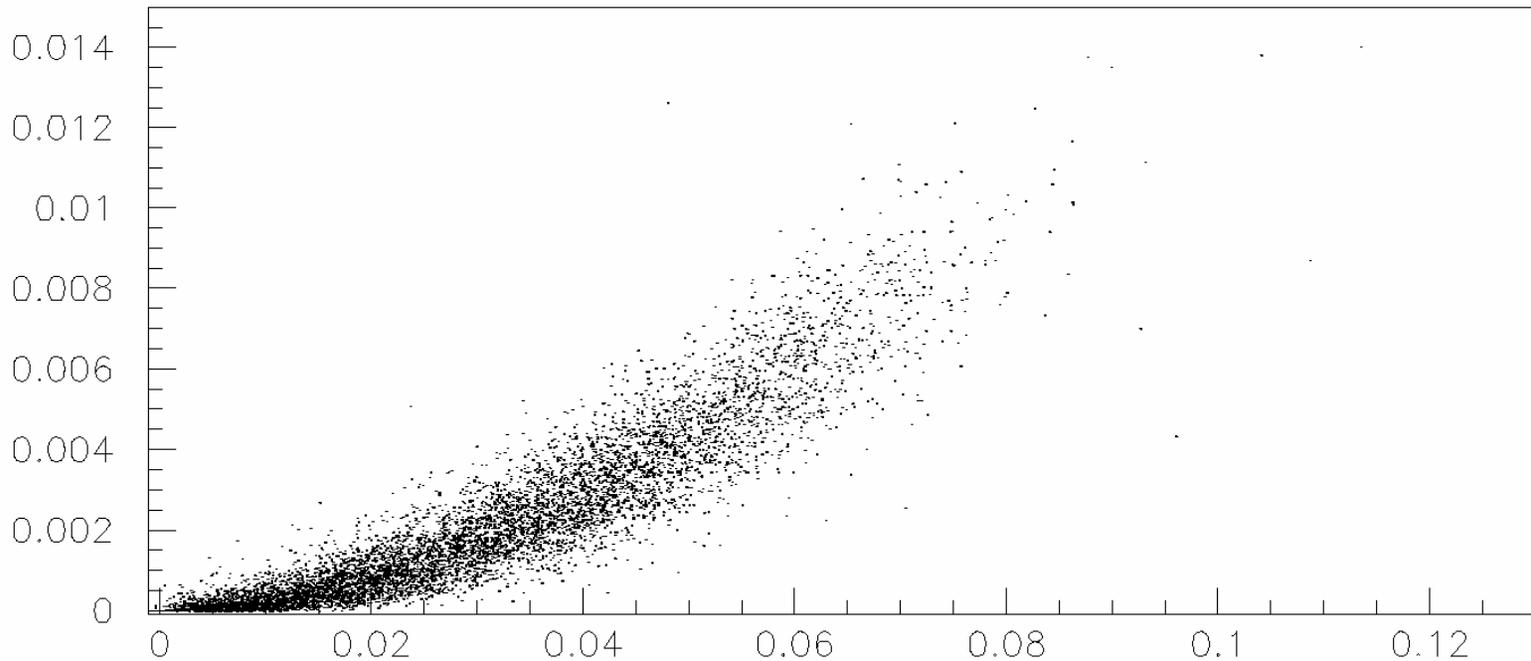


Figure taken from detector review document, January 1999

TWIST-TEC Projected Performance

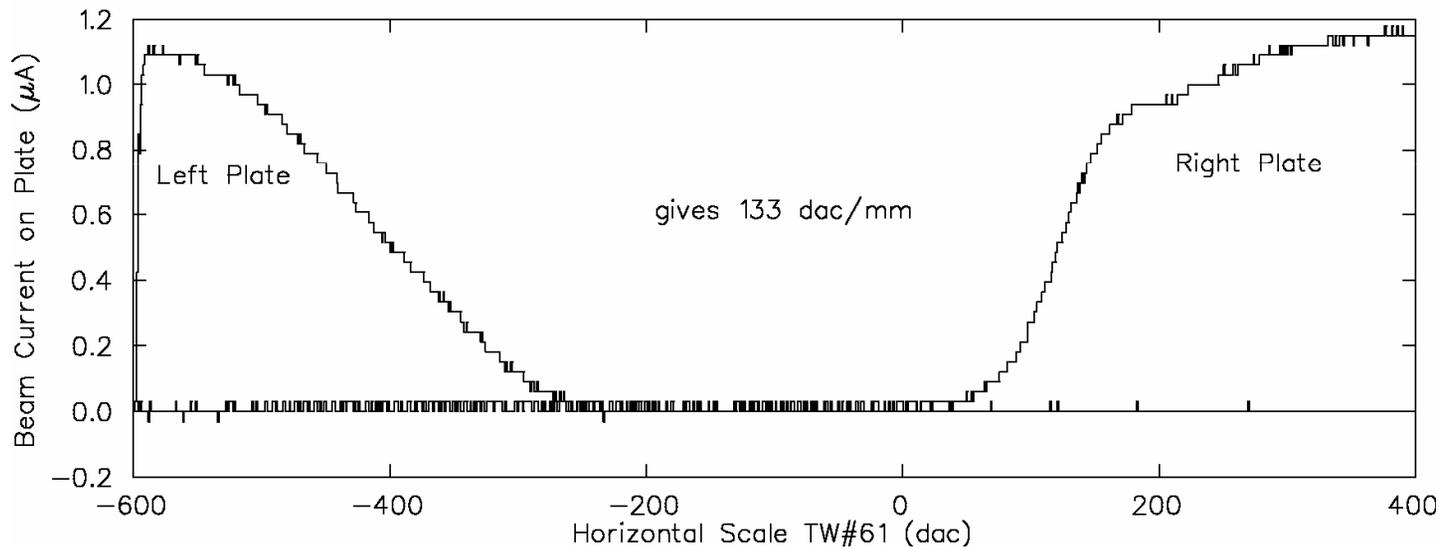
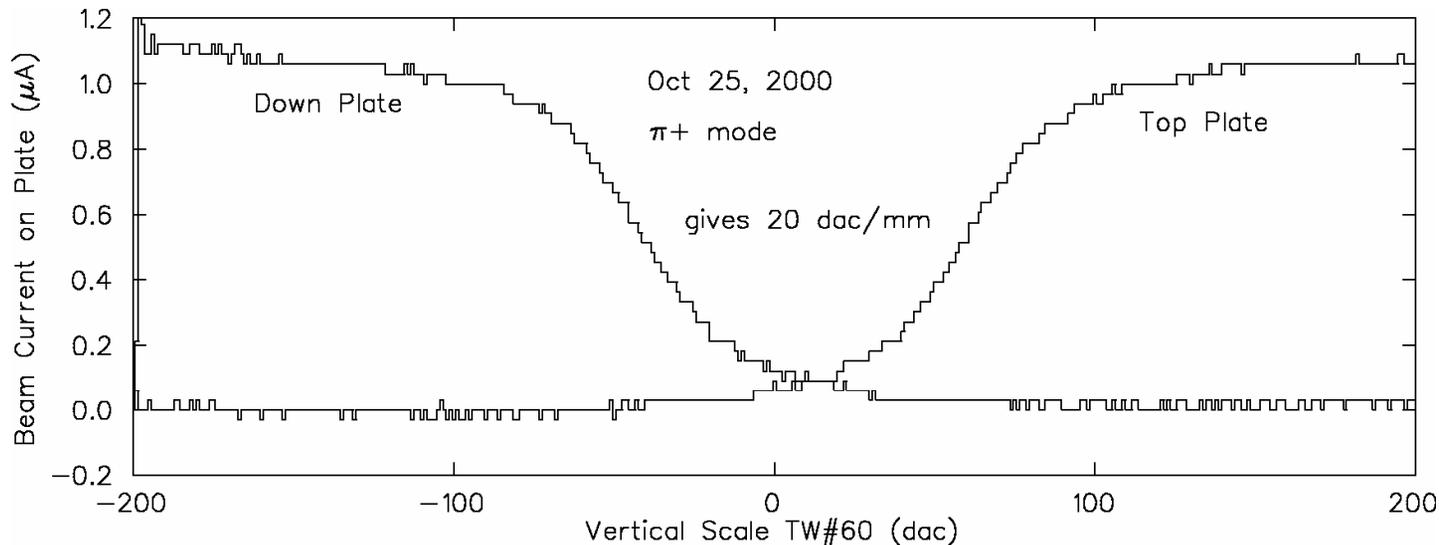
Effective Depolarization vs. TEC Tracking Angle



Sin(theta) (x-axis) and depolarization (y-axis)

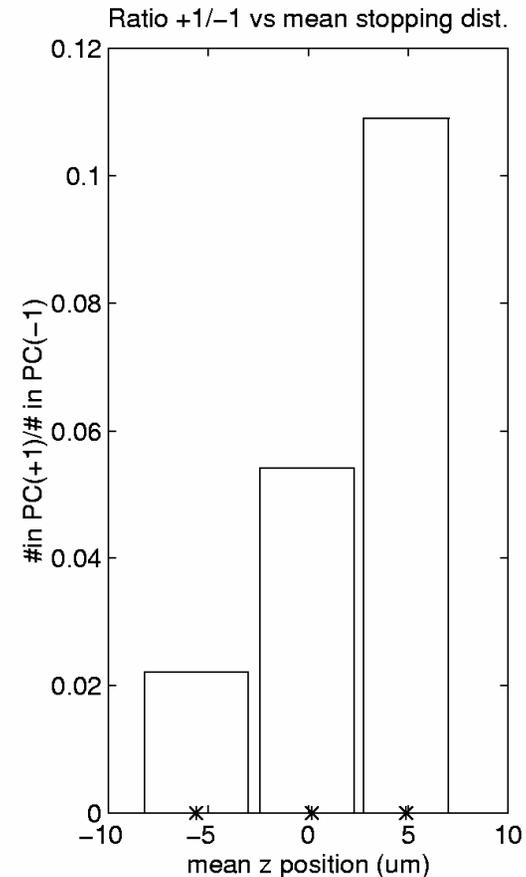
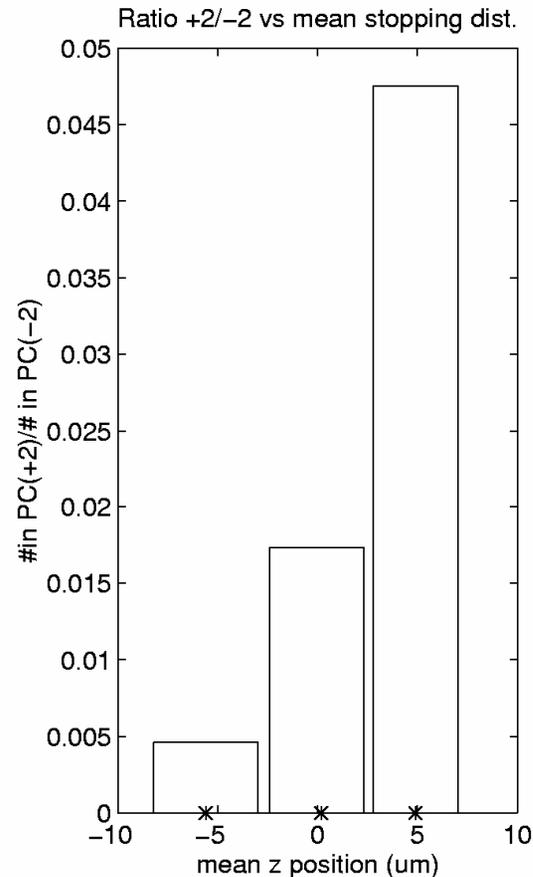
Correlation relies upon a highly convergent tune,
focused at the peak in the radial fringe field

TWIST-Proton Beam Monitor



TWIST - stopping distributions

Signal ratios in the target PC's can be used to monitor the stopping distribution

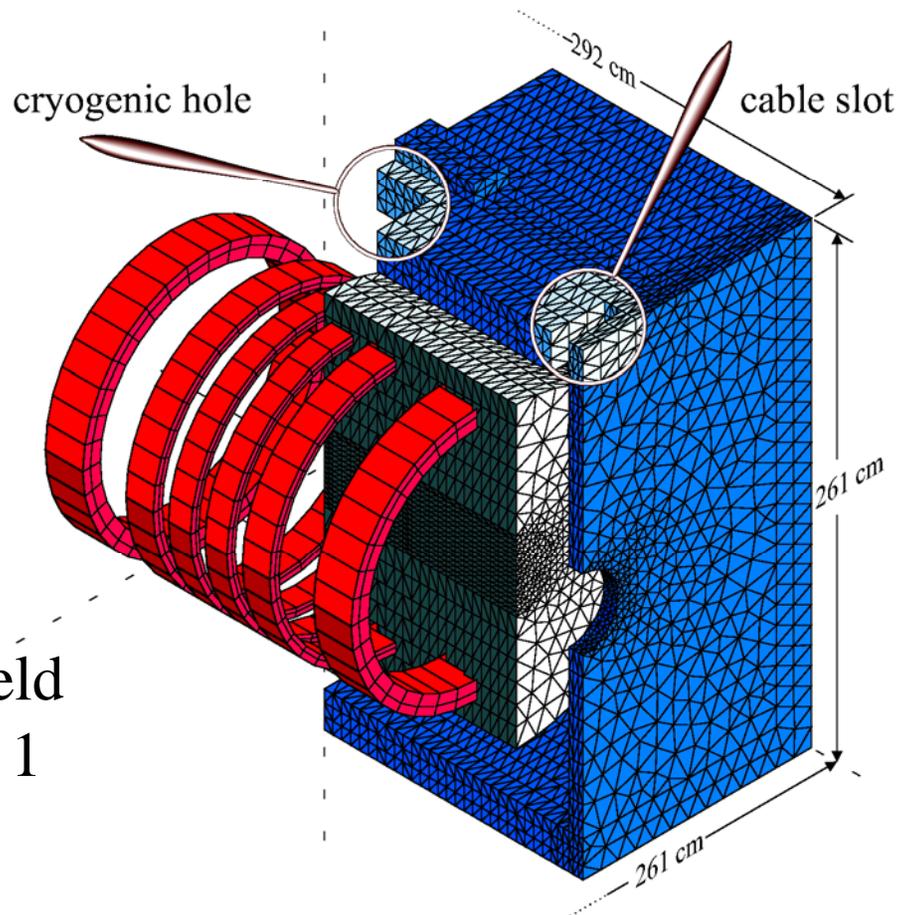


TWIST – Field Calculations

¼ yoke and all coils shown

- steel
- volume of reduced potential
- coils

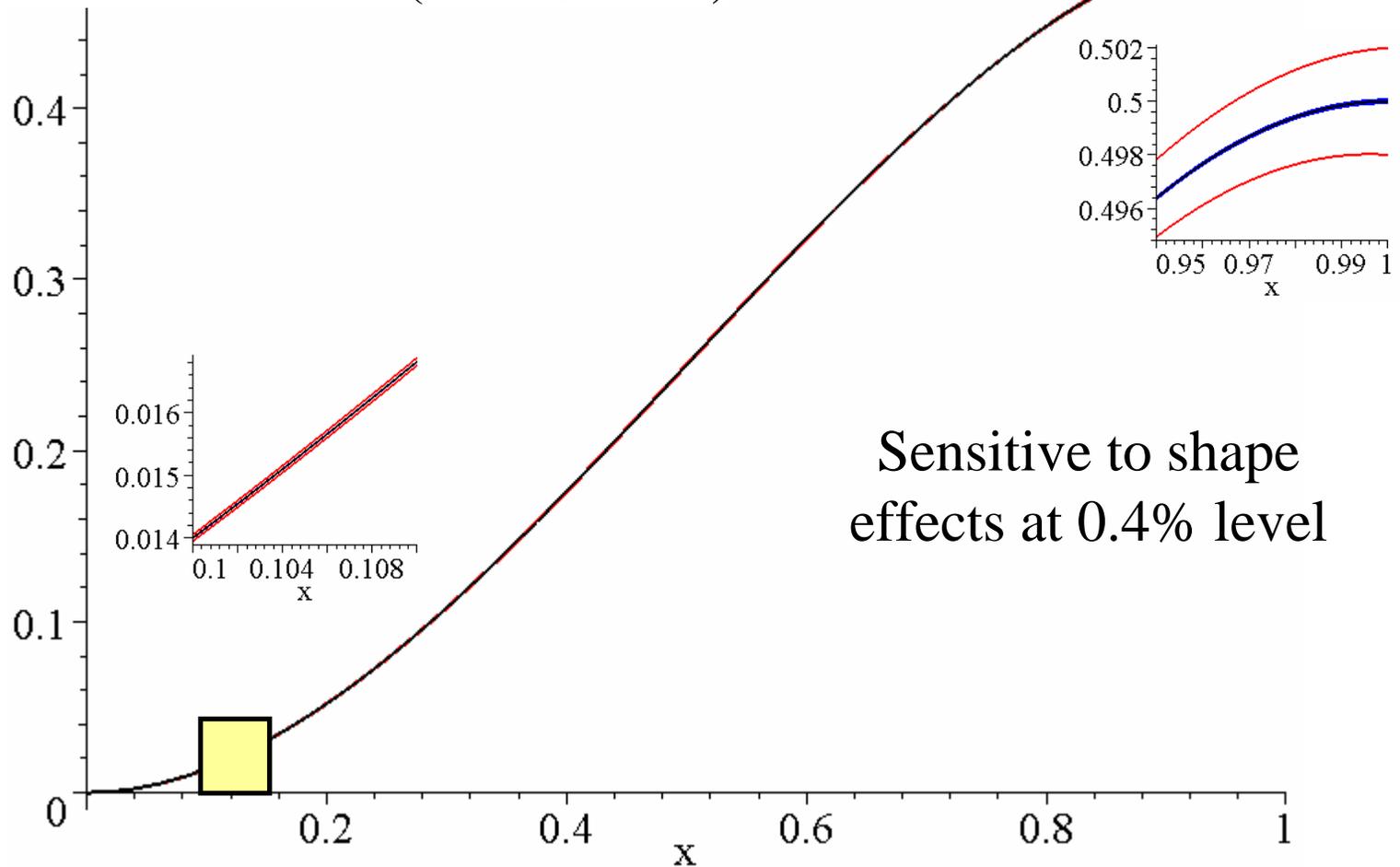
Dennis Wright
with OPERA-3d



Anticipated field
uniformity to 1
part in 10^4

Unpolarized distribution in x

$$[Forward + Backward] \propto 2x^2 \left(3(1-x) + \frac{2}{3} \rho(4x-3) \right)$$



Michel Parameters - defined

$$\begin{aligned} \rho \equiv & \frac{3}{4} \left[|g_{LL}^V|^2 + |g_{RR}^V|^2 + |g_{LR}^T|^2 + |g_{RL}^T|^2 \right] \\ & + \frac{3}{16} \left[|g_{LL}^S|^2 + |g_{RR}^S|^2 + |g_{LR}^S|^2 + |g_{RL}^S|^2 \right] \\ & - \frac{3}{4} \left[\text{Re}(g_{LR}^S g_{LR}^{T*}) + \text{Re}(g_{RL}^S g_{RL}^{T*}) \right] \end{aligned}$$

$$\begin{aligned} \xi \equiv & |g_{LL}^V|^2 + 3|g_{LR}^V|^2 - 3|g_{RL}^V|^2 - |g_{RR}^V|^2 + 5|g_{LR}^T|^2 \\ & - 5|g_{RL}^T|^2 + \frac{1}{4}|g_{LL}^S|^2 - \frac{1}{4}|g_{LR}^S|^2 + \frac{1}{4}|g_{RL}^S|^2 - \frac{1}{4}|g_{RR}^S|^2 \\ & + 4\text{Re}(g_{LR}^S g_{LR}^{T*}) - 4\text{Re}(g_{RL}^S g_{RL}^{T*}) \end{aligned}$$

$$\begin{aligned} \xi\delta \equiv & \frac{3}{4} \left[|g_{LL}^V|^2 - |g_{RR}^V|^2 - |g_{LR}^T|^2 + |g_{RL}^T|^2 \right] \\ & + \frac{3}{16} \left[|g_{LL}^S|^2 - |g_{RR}^S|^2 - |g_{LR}^S|^2 + |g_{RL}^S|^2 \right] \\ & - \frac{3}{4} \left[\text{Re}(g_{LR}^S g_{LR}^{T*}) - \text{Re}(g_{RL}^S g_{RL}^{T*}) \right] \end{aligned}$$